

Second-Order Linear Dynamic Systems

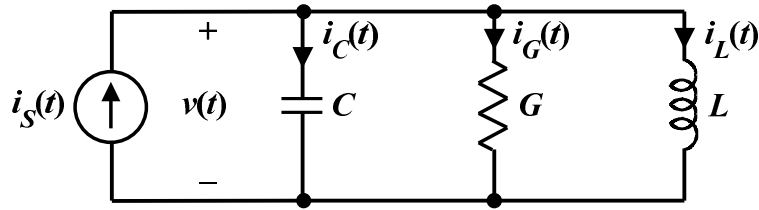
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Second-order linear dynamic systems are described by equations of the form:

$$\frac{d^2y(t)}{dt^2} + 2\zeta\omega_n \frac{dy}{dt} + \omega_n^2 y(t) = A\omega_n^2 z(t) \quad (1.1)$$

where $y(t)$ is the system response, or output, and $z(t)$ is the forcing function, or input. The symbols adopted here are a commonly used engineering notation, regardless of the field of concern. ζ is called the damping ratio, A is the *DC* or *static* gain, and ω_n is the natural frequency of the system. Several examples are given below.

Parallel RLC Circuit



Using the elementary component i - v relationships, we write:

$$v(t) = L \frac{di_L(t)}{dt} \quad (1.2)$$

$$i_G(t) = Gv(t) = G \left[L \frac{di_L(t)}{dt} \right] = GL \frac{di_L(t)}{dt} \quad (1.3)$$

$$i_C(t) = C \frac{dv(t)}{dt} = C \frac{d}{dt} \left[L \frac{di_L(t)}{dt} \right] = CL \frac{d^2i_L(t)}{dt^2} \quad (1.4)$$

Upon applying Kirchhoff's Current Law

$$i_C(t) + i_G(t) + i_L(t) = i_s(t) \quad (1.5)$$

we see that this circuit can be described by the second-order linear ordinary differential equation:

$$CL \frac{d^2i_L(t)}{dt^2} + GL \frac{di_L(t)}{dt} + i_L(t) = i_s(t) \quad (1.6)$$

or

$$\frac{d^2i_L(t)}{dt^2} + \frac{G}{C} \frac{di_L(t)}{dt} + \frac{1}{LC} i_L(t) = \frac{1}{LC} i_s(t) \quad (1.7)$$

Comparing this result to equation (1.1), we see that

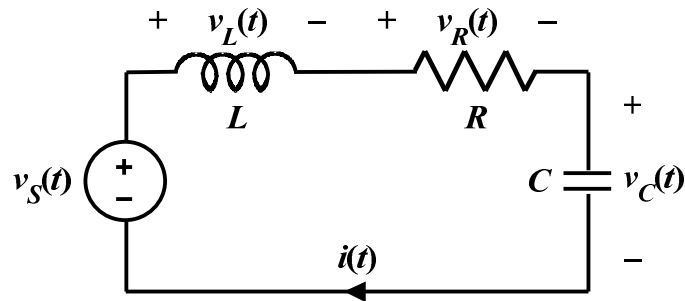
$$\omega_n^2 = \frac{1}{LC} \quad \Rightarrow \quad \omega_n = \frac{1}{\sqrt{LC}} \quad (1.8)$$

$$A\omega_n^2 = \frac{1}{LC} \quad \Rightarrow \quad A = 1 \quad (1.9)$$

and

$$2\zeta\omega_n = \frac{G}{C} \quad \Rightarrow \quad \zeta = \frac{G}{2\omega_n C} = \frac{G}{2} \sqrt{\frac{L}{C}} \quad (1.10)$$

Series RLC Circuit



Using the elementary component i - v relationships, we write:

$$i(t) = C \frac{dv_C(t)}{dt} \quad (1.11)$$

$$v_R(t) = Ri(t) = R \left[C \frac{dv_C(t)}{dt} \right] = RC \frac{dv_C(t)}{dt} \quad (1.12)$$

$$v_L(t) = L \frac{di(t)}{dt} = L \frac{d}{dt} \left[C \frac{dv_C(t)}{dt} \right] = LC \frac{d^2v_C(t)}{dt^2} \quad (1.13)$$

Upon applying Kirchhoff's Voltage Law

$$v_L(t) + v_R(t) + v_C(t) = v_S(t) \quad (1.14)$$

we see that this circuit can be described by the second-order linear ordinary differential equation:

$$LC \frac{d^2v_C(t)}{dt^2} + RC \frac{dv_C(t)}{dt} + v_C(t) = v_S(t) \quad (1.15)$$

or

$$\frac{d^2v_C(t)}{dt^2} + \frac{R}{L} \frac{dv_C(t)}{dt} + \frac{1}{LC} v_C(t) = \frac{1}{LC} v_S(t) \quad (1.16)$$

Comparing this result to equation (1.1), we see that

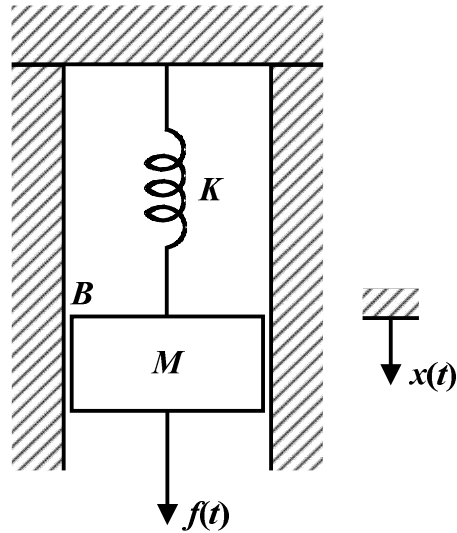
$$\omega_n^2 = \frac{1}{LC} \quad \Rightarrow \quad \omega_n = \frac{1}{\sqrt{LC}} \quad (1.17)$$

$$A\omega_n^2 = \frac{1}{LC} \quad \Rightarrow \quad A = 1 \quad (1.18)$$

and

$$2\zeta\omega_n = \frac{R}{L} \quad \Rightarrow \quad \zeta = \frac{R}{2\omega_n L} = \frac{R}{2} \sqrt{\frac{C}{L}} \quad (1.19)$$

Translational Mechanical System



When $f(t)$ is applied, friction and the spring will resist any motion so that, according to Newton's Second Law of Motion,

$$f(t) - Bv(t) - Kx(t) = Ma(t) \quad (1.20)$$

where $v(t) = \frac{dx(t)}{dt}$ is the velocity, and $a(t) = \frac{dv(t)}{dt} = \frac{d}{dt} \left[\frac{dx(t)}{dt} \right] = \frac{d^2x(t)}{dt^2}$ is the acceleration.

Substituting these into equation (1.20) yields the second-order linear ordinary differential equation:

$$f(t) - B \frac{dx(t)}{dt} - Kx(t) = M \frac{d^2x(t)}{dt^2} \quad (1.21)$$

or

$$\frac{d^2x(t)}{dt^2} + \frac{B}{M} \frac{dx(t)}{dt} + \frac{K}{M} x(t) = \frac{1}{M} f(t) \quad (1.22)$$

verifying that Newton's Second Law of Motion is clearly a mechanical equivalent to Kirchhoff's Laws for electrical circuits. Comparing this result to equation (1.1), we see that

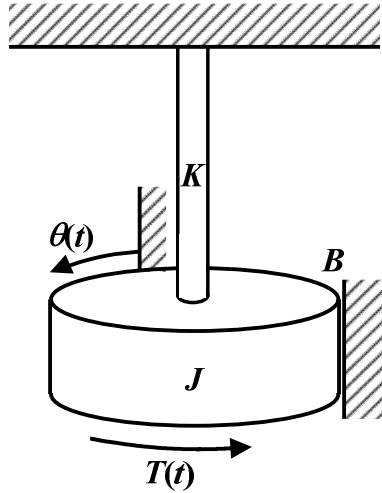
$$\omega_n^2 = \frac{K}{M} \Rightarrow \omega_n = \sqrt{\frac{K}{M}} \quad (1.23)$$

$$A\omega_n^2 = \frac{1}{M} \Rightarrow A = \frac{1}{K} \quad (1.24)$$

and

$$2\zeta\omega_n = \frac{B}{M} \Rightarrow \zeta = \frac{B}{2\omega_n M} = \frac{B}{2\sqrt{KM}} \quad (1.25)$$

Rotational Mechanical System



When $T(t)$ is applied, friction and the spring will resist any motion so that, according to Newton's Second Law of Motion,

$$T(t) - B\omega(t) - K\theta(t) = J\alpha(t) \quad (1.26)$$

where $\omega(t) = \frac{d\theta(t)}{dt}$ is the angular velocity, and $\alpha(t) = \frac{d\omega(t)}{dt} = \frac{d}{dt} \left[\frac{d\theta(t)}{dt} \right] = \frac{d^2\theta(t)}{dt^2}$ is the angular acceleration. Substituting these into equation (1.26) yields the second-order linear ordinary differential equation:

$$T(t) - B \frac{d\theta(t)}{dt} - K\theta(t) = J \frac{d^2\theta(t)}{dt^2} \quad (1.27)$$

or

$$\frac{d^2\theta(t)}{dt^2} + \frac{B}{J} \frac{d\theta(t)}{dt} + \frac{K}{J} \theta(t) = \frac{1}{J} T(t) \quad (1.28)$$

which again is clearly analogous to Kirchhoff's Laws for electrical circuits. Comparing this result to equation (1.1), we see that

$$\omega_n^2 = \frac{K}{J} \Rightarrow \omega_n = \sqrt{\frac{K}{J}} \quad (1.29)$$

$$A\omega_n^2 = \frac{1}{J} \Rightarrow A = \frac{1}{K} \quad (1.30)$$

and

$$2\zeta\omega_n = \frac{B}{J} \Rightarrow \zeta = \frac{B}{2\omega_n J} = \frac{B}{2\sqrt{KJ}} \quad (1.31)$$

All of the examples considered here yield equations that are of the form of equation (1.1). Note that, if $z(t) = 0$, the differential equation is said to be homogeneous, and the system response under that condition is called the *natural* response. If $z(t) \neq 0$, the differential equation is said to be non-homogeneous, and the *complete* response of the system with the forcing function applied is a combination of the *natural* response and additional term(s) called the *forced* response. Sometimes, these are called, respectively, the *complimentary* response and the *particular* response.

Zero-Input (Unforced) Systems

Consider the zero-input (homogeneous) form of equation (1.1):

$$\frac{d^2 y(t)}{dt^2} + 2\zeta\omega_n \frac{dy(t)}{dt} + \omega_n^2 y(t) = 0 \quad (1.32)$$

If we assume that the natural response of the system is exponential, i.e., $y(t) = \beta e^{rt}$, then

$$r^2 \beta e^{rt} + 2\zeta\omega_n r \beta e^{rt} + \omega_n^2 \beta e^{rt} = 0 \quad (1.33)$$

or

$$(r^2 + 2\zeta\omega_n r + \omega_n^2) \beta e^{rt} = 0 \quad (1.34)$$

which means that

$$r^2 + 2\zeta\omega_n r + \omega_n^2 = 0 \quad (1.35)$$

Equation (1.35) is called the characteristic equation of the system, and it has roots given by:

$$\begin{aligned} r_{1,2} &= \frac{-2\zeta\omega_n \pm \sqrt{(2\zeta\omega_n)^2 - 4\omega_n^2}}{2} \\ &= -\zeta\omega_n \pm \omega_n \sqrt{\zeta^2 - 1} \\ &= \left(-\zeta \pm \sqrt{\zeta^2 - 1}\right) \omega_n \end{aligned} \quad (1.36)$$

From this, we will see that there are four distinctly different forms of the solution to equation (1.32), depending on the value of ζ with respect to the number 1.

Case 1

If $\zeta > 1$, then $\zeta^2 - 1 > 0$, and there will be two distinct negative real roots, $r_1 = (-\zeta + \sqrt{\zeta^2 - 1})\omega_n$ and $r_2 = (-\zeta - \sqrt{\zeta^2 - 1})\omega_n$. In this case, the system is said to be **overdamped**, and because there are **two** roots to the characteristic equation, $y(t)$ will have **two** exponential components:

$$y(t) = \beta_1 e^{r_1 t} + \beta_2 e^{r_2 t} \quad (1.37)$$

To determine the values of β_1 and β_2 note that

$$\dot{y}(t) = \beta_1 r_1 e^{r_1 t} + \beta_2 r_2 e^{r_2 t} \quad (1.38)$$

Evaluating equations (1.37) and (1.38) at $t = 0$, we have

$$\beta_1 + \beta_2 = y(0) \quad (1.39)$$

and

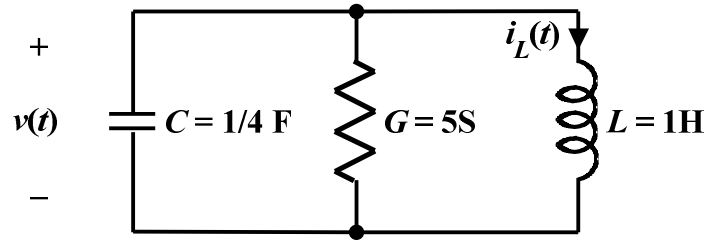
$$r_1 \beta_1 + r_2 \beta_2 = \dot{y}(0) \quad (1.40)$$

These two simultaneous equations can be used to evaluate β_1 and β_2 using Cramer's Rule as follows:

$$\beta_1 = \frac{\begin{vmatrix} y(0) & 1 \\ \dot{y}(0) & r_2 \end{vmatrix}}{\begin{vmatrix} 1 & 1 \\ r_1 & r_2 \end{vmatrix}} = \frac{r_2 y(0) - \dot{y}(0)}{r_2 - r_1} \quad (1.41)$$

$$\beta_2 = \frac{\begin{vmatrix} 1 & y(0) \\ r_1 & \dot{y}(0) \end{vmatrix}}{\begin{vmatrix} 1 & 1 \\ r_1 & r_2 \end{vmatrix}} = \frac{\dot{y}(0) - r_1 y(0)}{r_2 - r_1} \quad (1.42)$$

Example 1.1



As shown by equation (1.7), this parallel RLC circuit can be described by the equation

$$\frac{d^2 i_L}{dt^2} + 20 \frac{di_L}{dt} + 4i_L = 0 \quad (1.43)$$

Hence, the characteristic equation is

$$r^2 + 20r + 4 = 0 \quad (1.44)$$

and

$$\omega_n = \frac{1}{\sqrt{1(1/4)}} = 2 \quad (1.45)$$

$$\zeta = \frac{5}{2} \sqrt{\frac{1}{(1/4)}} = 5 \quad (1.46)$$

This is an overdamped system, with

$$r_1 = (-5 + \sqrt{25-1})2 \approx -0.202 \quad (1.47)$$

$$r_2 = (-5 - \sqrt{25-1})2 \approx -19.798 \quad (1.48)$$

Suppose now that $i_L(0) = 0$ and $v(0) = 1$. Then, $v(t) = L \frac{di_L(t)}{dt}$, when evaluated at $t = 0$,

yields $\left. \frac{di_L}{dt} \right|_{t=0} = \frac{1}{L} v(0) = 1$, and

$$\beta_1 \approx \frac{(-19.798)(0) - 1}{-19.798 - (-0.202)} \approx \frac{-1}{-19.596} \approx 0.051 \quad (1.49)$$

$$\beta_2 \approx \frac{1 - (-0.202)(0)}{-19.798 - (-0.202)} \approx \frac{1}{-19.596} \approx -0.051 \quad (1.50)$$

Hence,

$$i_L(t) \approx 0.051e^{-0.202t} - 0.051e^{-19.798t} \text{ A for } t > 0 \quad (1.51)$$

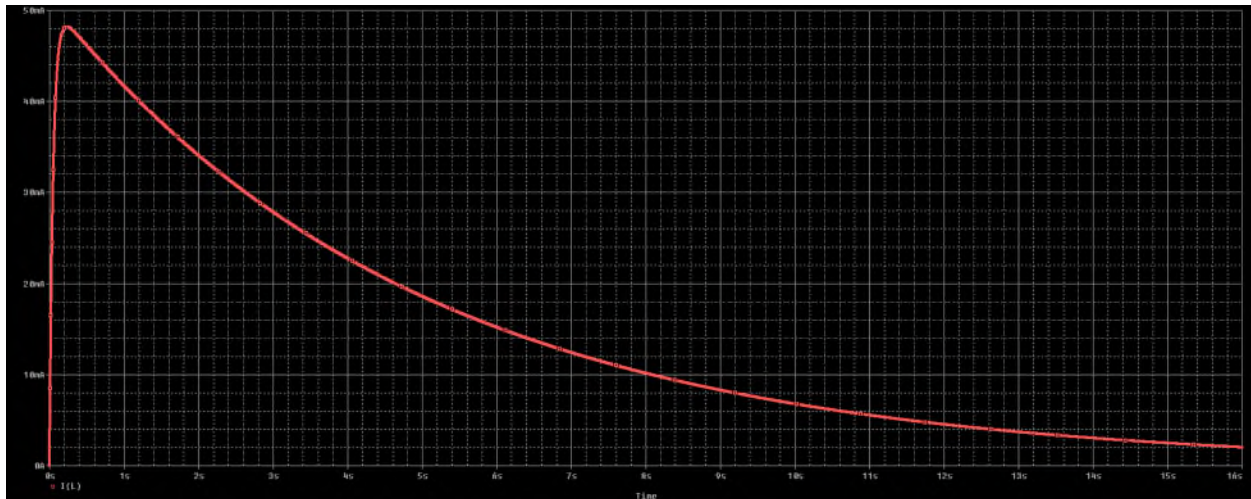
To see what this looks like, we can simulate the circuit with PSpice as follows:

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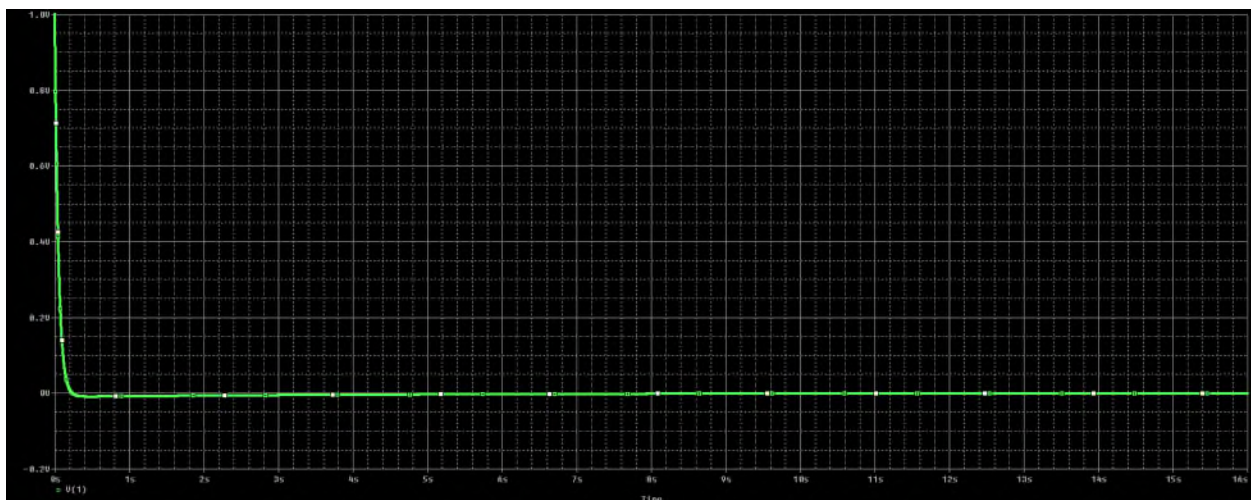
Example 1.1
C      1      0      {1/4}      IC=1
G      1      0      1      0      5
L      1      0      1      IC=0
.TRAN  1      16      0      1m      UIC
.PROBE
.END

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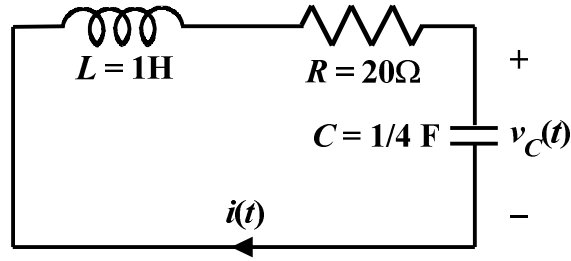
The inductor current is:



and the capacitor voltage is:



Example 1.2



As shown by equation (1.16), this series RLC circuit can be described by the equation

$$\frac{d^2 v_C}{dt^2} + 20 \frac{dv_C}{dt} + 4v_C = 0 \quad (1.52)$$

Hence, the characteristic equation is

$$r^2 + 20r + 4 = 0 \quad (1.53)$$

and

$$\omega_n = \frac{1}{\sqrt{1(1/4)}} = 2 \quad (1.54)$$

$$\zeta = \frac{20}{2} \sqrt{\frac{(1/4)}{1}} = 5 \quad (1.55)$$

This is an overdamped system, with

$$r_1 = (-5 + \sqrt{25-1})2 \approx -0.202 \quad (1.56)$$

$$r_2 = (-5 - \sqrt{25-1})2 \approx -19.798 \quad (1.57)$$

Suppose now that $v_C(0) = 0$ and $i(0) = 1$. Then, $i(t) = C \frac{dv_C(t)}{dt}$, when evaluated at $t = 0$,

yields $\left. \frac{dv_C(t)}{dt} \right|_{t=0} = \frac{1}{C} i(0) = 4$, and

$$\beta_1 \approx \frac{(-19.798)(0) - 4}{-19.798 - (-0.202)} \approx 0.204 \quad (1.58)$$

$$\beta_2 \approx \frac{4 - (-0.202)(0)}{-19.798 - (-0.202)} \approx -0.204 \quad (1.59)$$

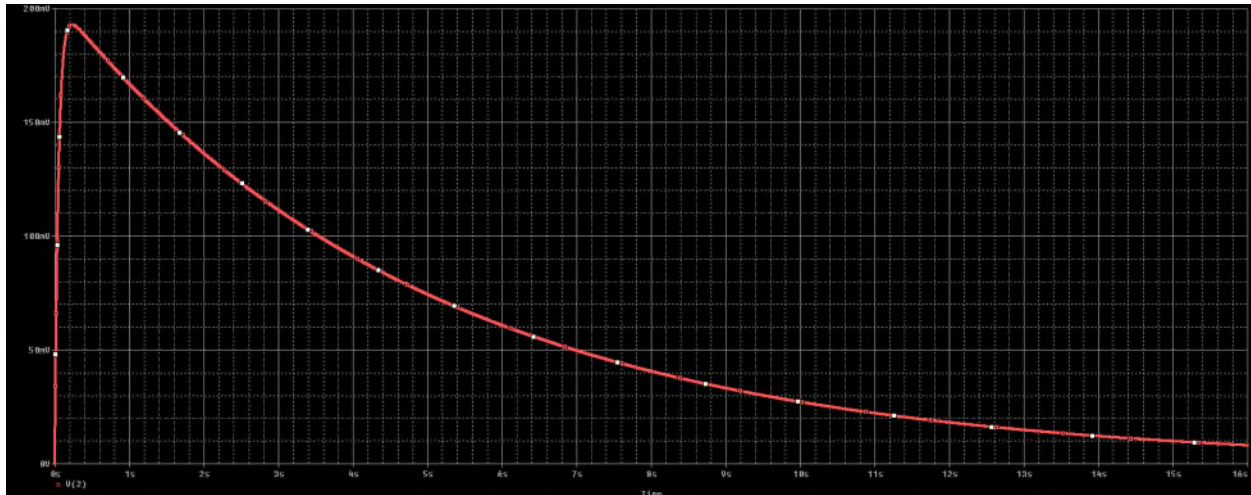
Hence,

$$v_C(t) \approx 0.204e^{-0.202t} - 0.204e^{-19.798t} \text{ V for } t > 0 \quad (1.60)$$

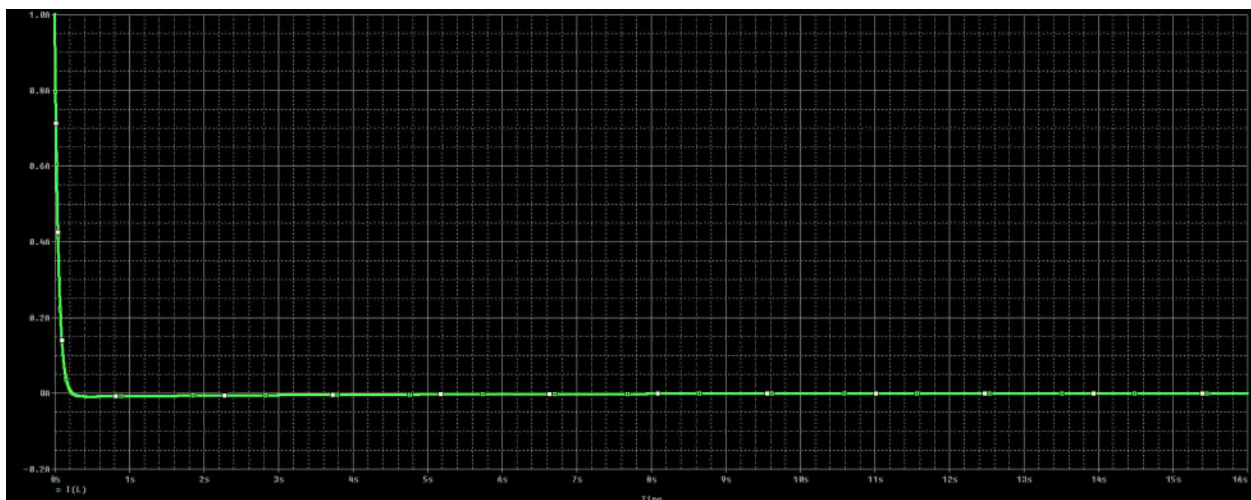
To see what this looks like, we can simulate the circuit with PSpice as follows:

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Example 1.2
L      0      1      1      IC=1
R      1      2      20
C      2      0      {1/4}      IC=0
.TRAN  1      16      0      1m      UIC
.PROBE
.END
```

The capacitor voltage is:



and the inductor current is:



Case 2

If $0 < \zeta < 1$, then $\zeta^2 - 1 < 0$, and there will be two complex conjugate roots,

$r_1 = \left(-\zeta + j\sqrt{1-\zeta^2}\right)\omega_n = -\zeta\omega_n + j\omega_d$ and $r_2 = \left(-\zeta - j\sqrt{1-\zeta^2}\right)\omega_n = -\zeta\omega_n - j\omega_d$. In this case,

the system is said to be **underdamped**., and the quantity $\omega_d = \omega_n\sqrt{1-\zeta^2}$ is called the *damped* or *ringing* frequency.

As in Case 1, because there are **two** distinct roots to the characteristic equation, $y(t)$ has **two** exponential components:

$$\begin{aligned} y(t) &= \beta_1 e^{(-\zeta\omega_n + j\omega_d)t} + \beta_2 e^{(-\zeta\omega_n - j\omega_d)t} \\ &= e^{-\zeta\omega_n t} \left(\beta_1 e^{j\omega_d t} + \beta_2 e^{-j\omega_d t} \right) \end{aligned} \quad (1.61)$$

However, it is usually preferred to use Euler's identity

$$e^{\pm j\theta} = \cos \theta \pm j \sin \theta \quad (1.62)$$

to express $y(t)$ in the alternate form

$$\begin{aligned} y(t) &= e^{-\zeta\omega_n t} \left[\beta_1 (\cos \omega_d t + j \sin \omega_d t) + \beta_2 (\cos \omega_d t - j \sin \omega_d t) \right] \\ &= e^{-\zeta\omega_n t} \left[(\beta_1 + \beta_2) \cos \omega_d t + j(\beta_1 - \beta_2) \sin \omega_d t \right] \\ &= e^{-\zeta\omega_n t} \left[B_1 \cos \omega_d t + B_2 \sin \omega_d t \right] \end{aligned} \quad (1.63)$$

where $B_1 = \beta_1 + \beta_2$ and $B_2 = j(\beta_1 - \beta_2)$.

To determine the values of B_1 and B_2 note that

$$\dot{y}(t) = -\zeta\omega_n e^{-\zeta\omega_n t} \left[B_1 \cos \omega_d t + B_2 \sin \omega_d t \right] + e^{-\zeta\omega_n t} \left[-B_1 \omega_d \sin \omega_d t + B_2 \omega_d \cos \omega_d t \right] \quad (1.64)$$

Evaluating equations (1.63) and (1.64) at $t = 0$, we have

$$B_1 = y(0) \quad (1.65)$$

and

$$-\zeta\omega_n B_1 + B_2 \omega_d = \dot{y}(0) \quad (1.66)$$

Thus,

$$B_2 = \frac{\dot{y}(0) + \zeta\omega_n B_1}{\omega_d} = \frac{\dot{y}(0) + \zeta\omega_n y(0)}{\omega_d} \quad (1.67)$$

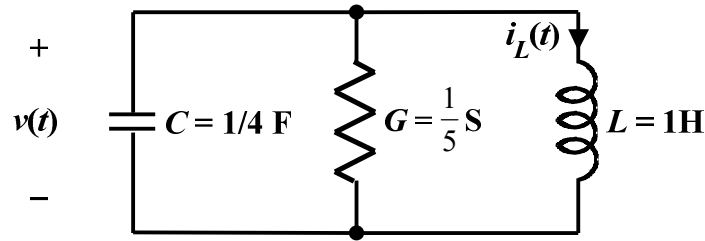
Alternately, note that

$$B_1 \cos \omega_d t + B_2 \sin \omega_d t = B_3 \cos(\omega_d t - \phi) \quad (1.68)$$

where $B_3 = \sqrt{B_1^2 + B_2^2}$ and $\phi = \tan^{-1}\left(\frac{B_2}{B_1}\right)$, so that $y(t)$ can be written in a slightly more compact form as

$$y(t) = B_3 e^{-\zeta\omega_n t} \cos(\omega_d t - \phi) \quad (1.69)$$

Example 2.1



As shown by equation (1.7), this parallel RLC circuit can be described by the equation

$$\frac{d^2 i_L}{dt^2} + \frac{4}{5} \frac{di_L}{dt} + 4i_L = 0 \quad (1.70)$$

Hence, the characteristic equation is

$$r^2 + \frac{4}{5}r + 4 = 0 \quad (1.71)$$

and

$$\omega_n = 2 \quad (1.72)$$

$$\zeta = \frac{1}{10} \sqrt{4} = 0.2 \quad (1.73)$$

$$\omega_d = 2\sqrt{1 - (0.2)^2} \approx 1.960 \quad (1.74)$$

This is an underdamped system, with

$$r_{1,2} \approx -0.400 + j1.960 \quad (1.75)$$

$$r_{1,2} \approx -0.400 - j1.960 \quad (1.76)$$

Suppose now that $i_L(0) = 0$ and $v(0) = 1$. Then, $v(t) = L \frac{di_L(t)}{dt}$, when evaluated at $t = 0$,

yields $\left. \frac{di_L(t)}{dt} \right|_{t=0} = \frac{1}{L} v(0) = 1$, and

$$B_1 = 0 \quad (1.77)$$

$$B_2 \approx \frac{1 + (0.2)(2)(0)}{1.960} \approx 0.510 \quad (1.78)$$

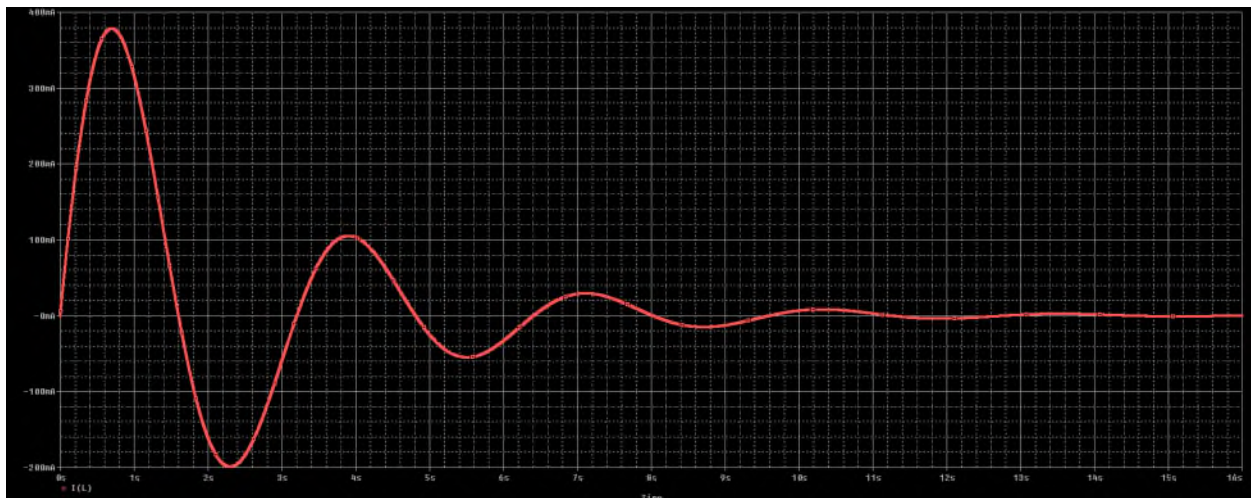
Hence,

$$i_L(t) \approx 0.510e^{-0.4t} \sin(1.960t) \text{ A for } t > 0 \quad (1.79)$$

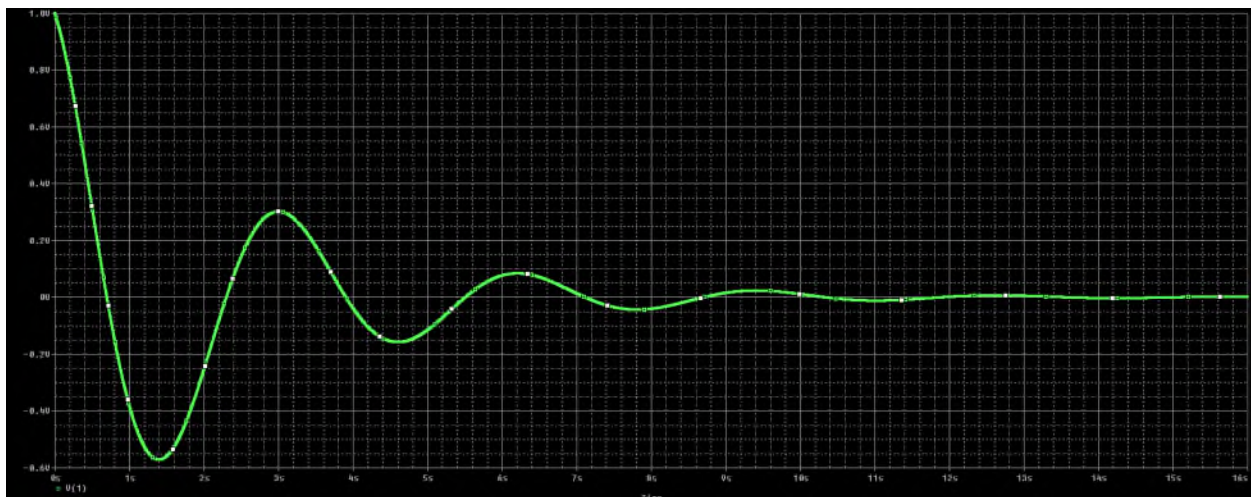
To see what this looks like, we can simulate the circuit with PSpice as follows:

```
Example 2.1
C 1 0 {1/4} IC=1
G 1 0 1 0 {1/5}
L 1 0 1 IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END
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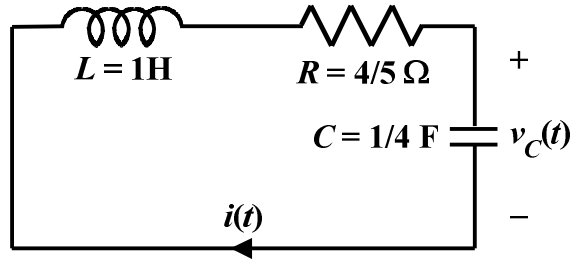
The inductor current is:



and the capacitor voltage is:



Example 2.2



As shown by equation (1.16), this series RLC circuit can be described by the equation

$$\frac{d^2 v_C}{dt^2} + \frac{4}{5} \frac{dv_C}{dt} + 4v_C = 0 \quad (1.80)$$

Hence, the characteristic equation is

$$r^2 + \frac{4}{5}r + 4 = 0 \quad (1.81)$$

and

$$\omega_n = 2 \quad (1.82)$$

$$\zeta = \frac{2}{5} \sqrt{\frac{1}{4}} = 0.2 \quad (1.83)$$

$$\omega_d = 2\sqrt{1 - (0.2)^2} \approx 1.960 \quad (1.84)$$

This is an underdamped system, with

$$r_1 = -0.400 + j1.960 \quad (1.85)$$

$$r_2 = -0.400 - j1.960 \quad (1.86)$$

Suppose now that $v_C(0) = 0$ and $i(0) = 1$. Then, $i(t) = C \frac{dv_C(t)}{dt}$, when evaluated at $t = 0$,

yields $\left. \frac{dv_C(t)}{dt} \right|_{t=0} = \frac{1}{C} i(0) = 4$, and

$$B_1 = 0 \quad (1.87)$$

$$B_2 \approx \frac{4 + (0.2)(2)(0)}{1.960} \approx 2.041 \quad (1.88)$$

Hence,

$$v_c(t) = 2.041e^{-0.4t} \sin(1.960t) \quad \forall t > 0 \quad (1.89)$$

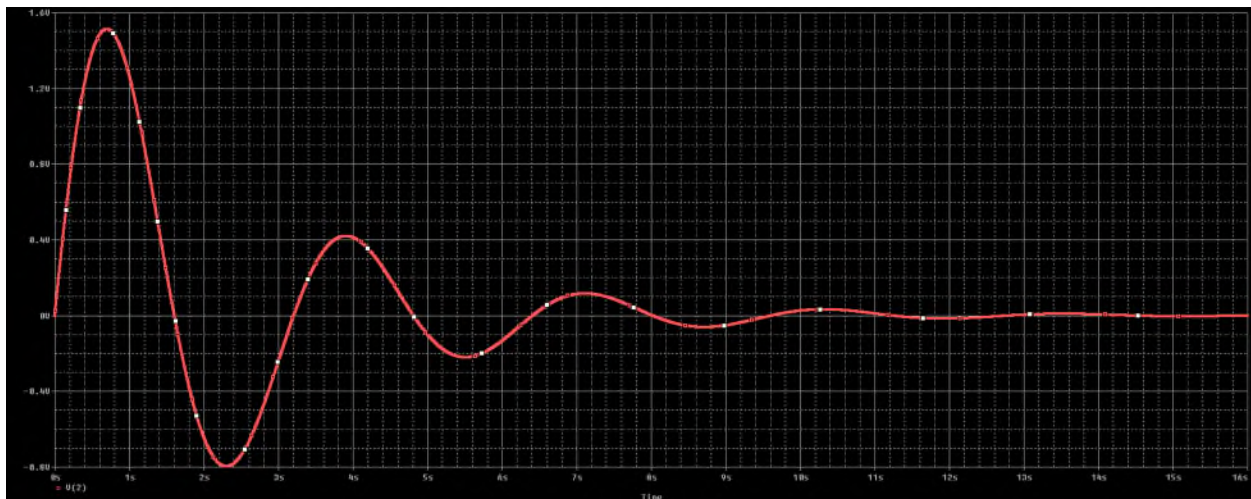
To see what this looks like, we can simulate the circuit with PSpice as follows:

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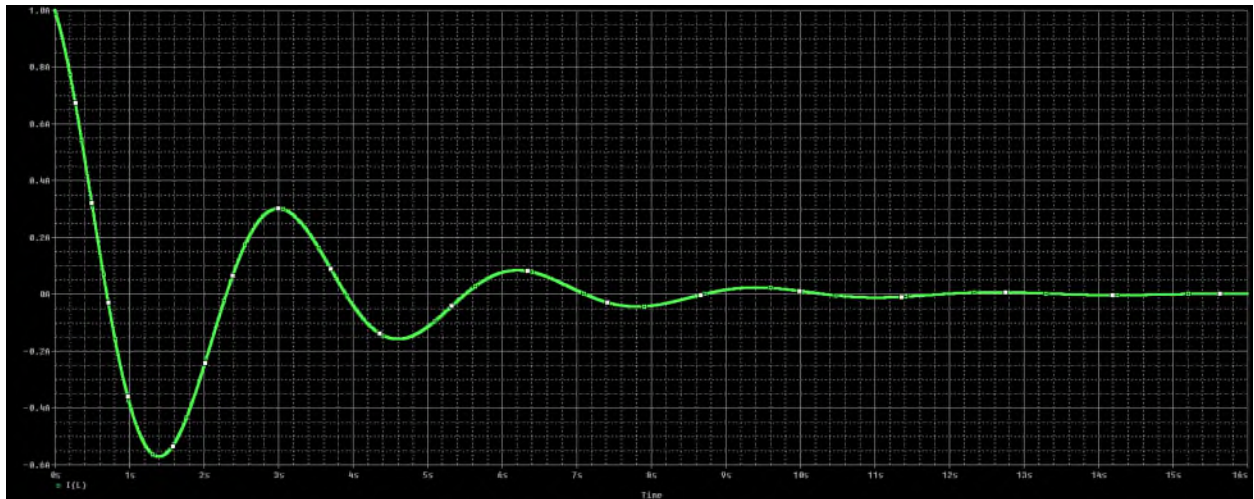
Example 2.2
L      0      1      1      IC=1
R      1      2      {4/5}
C      2      0      {1/4}      IC=0
.TRAN  1      16      0      1m      UIC
.PROBE
.END

```

The capacitor voltage is:



and the inductor current is:



Case 3

If $\zeta = 1$, then $\zeta^2 - 1 = 0$, and there will be two identical negative real roots, $r_1 = r_2 = -\omega_n$. In this case, the system is said to be *critically damped*. This case can be considered to be the “borderline” between overdamped and underdamped systems.

The general form of the solution is

$$y(t) = (\beta_1 + \beta_2 t) e^{-\omega_n t} \quad (1.90)$$

To determine the values of β_1 and β_2 note that

$$\dot{y}(t) = \beta_2 e^{-\omega_n t} - \omega_n (\beta_1 + \beta_2 t) e^{-\omega_n t} \quad (1.91)$$

Evaluating equations (1.90) and (1.91) at $t = 0$, we have

$$\beta_1 = y(0) \quad (1.92)$$

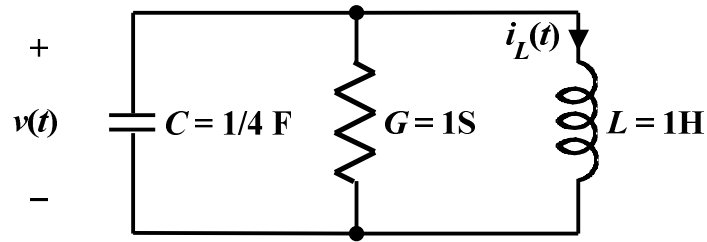
and

$$\beta_2 - \omega_n \beta_1 = \dot{y}(0) \quad (1.93)$$

Thus,

$$\beta_2 = \dot{y}(0) + \omega_n \beta_1 = \dot{y}(0) + \omega_n y(0) \quad (1.94)$$

Example 3.1



As shown by equation (1.7), this parallel circuit can be described by the equation

$$\frac{d^2 i_L}{dt^2} + 4 \frac{di_L}{dt} + 4i_L = 0 \quad (1.95)$$

Hence, the characteristic equation is

$$r^2 + 4r + 4 = 0 \quad (1.96)$$

and

$$\omega_n = 2 \quad (1.97)$$

$$\zeta = \frac{1}{2} \sqrt{4} = 1 \quad (1.98)$$

This is a critically damped system, with

$$r_1 = r_2 = -2 \quad (1.99)$$

Suppose now that $i_L(0) = 0$ and $v(0) = 1$. Then, $v(t) = L \frac{di_L(t)}{dt}$, when evaluated at $t = 0$,

yields $\left. \frac{di_L}{dt} \right|_{t=0} = \frac{1}{L} v(0) = 1$, and

$$\beta_1 = 0 \quad (1.100)$$

$$\beta_2 = 1 + (2)(0) = 1 \quad (1.101)$$

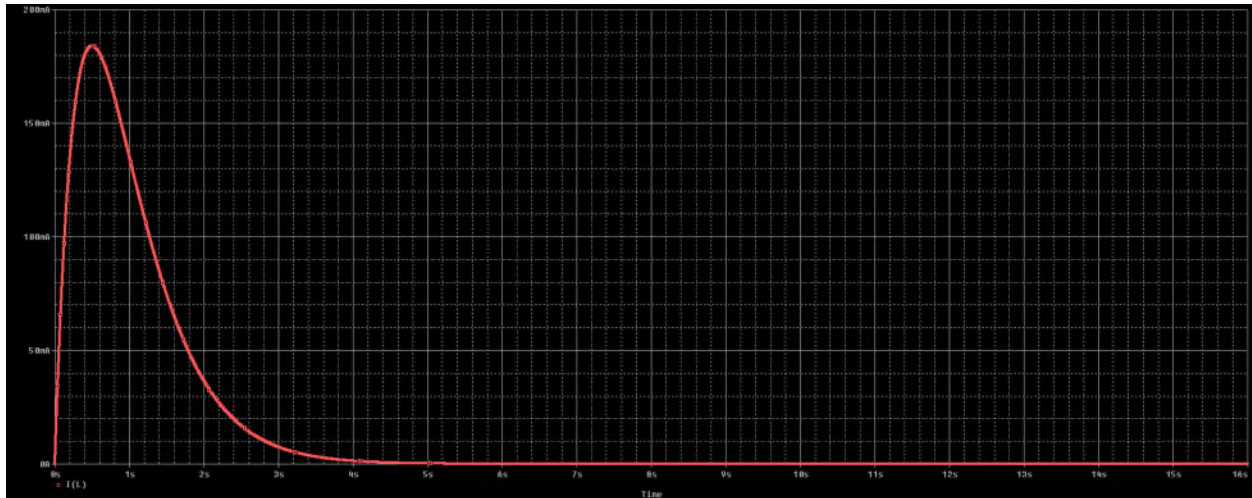
Hence,

$$i_L(t) = te^{-2t} \text{ A for } t > 0 \quad (1.102)$$

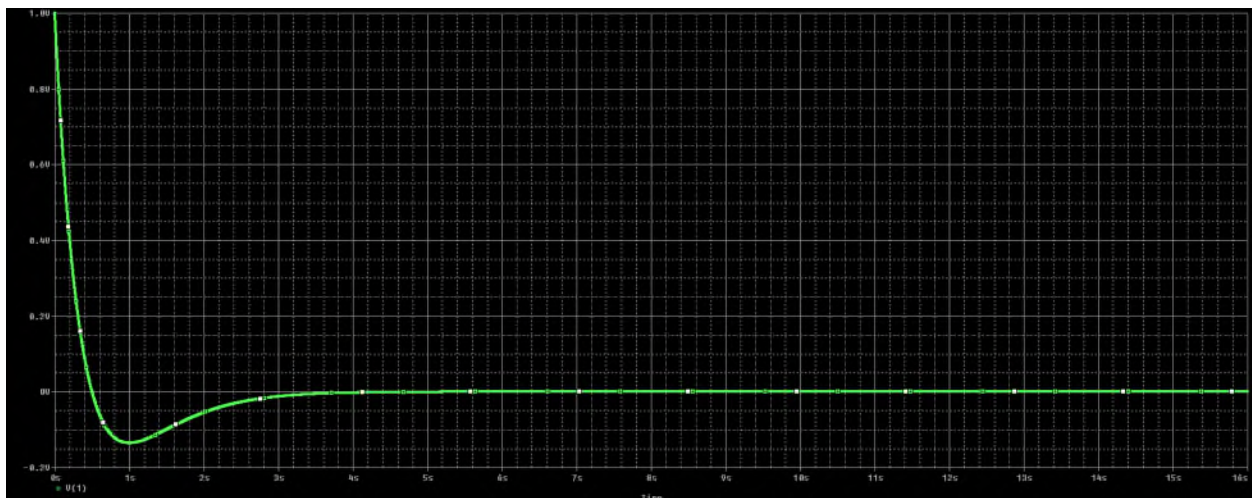
To see what this looks like, we can simulate the circuit with PSpice as follows:

```
Example 3.1
C 1 0 {1/4} IC=1
G 1 0 1 0 1
L 1 0 1 IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END
```

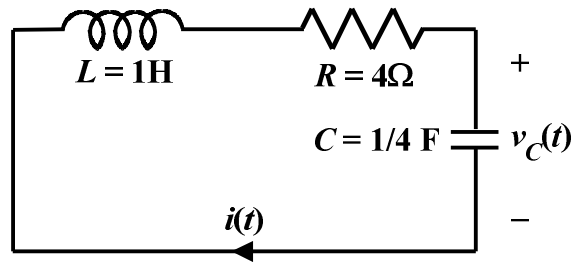
The inductor current is:



and the capacitor voltage is:



Example 3.2



As shown by equation (1.16), this series RLC circuit can be described by the equation

$$\frac{d^2v_C}{dt^2} + 4\frac{dv_C}{dt} + 4v_C = 0 \quad (1.103)$$

Hence, the characteristic equation is

$$r^2 + 4r + 4 = 0 \quad (1.104)$$

and

$$\omega_n = 2 \quad (1.105)$$

$$\zeta = \frac{4}{2}\sqrt{\frac{1}{4}} = 1 \quad (1.106)$$

This is a critically damped system, with

$$r_1 = r_2 = -2 \quad (1.107)$$

Suppose now that $v_C(0) = 0$ and $i(0) = 1$. Then, $i(t) = C\frac{dv_C(t)}{dt}$, when evaluated at $t = 0$,

yields $\left.\frac{dv_C(t)}{dt}\right|_{t=0} = \frac{1}{C}i(0) = 4$, and

$$\beta_1 = 0 \quad (1.108)$$

$$\beta_2 = 4 + (2)(0) = 4 \quad (1.109)$$

Hence,

$$v_C(t) = 4te^{-2t} \text{ V } t > 0 \quad (1.110)$$

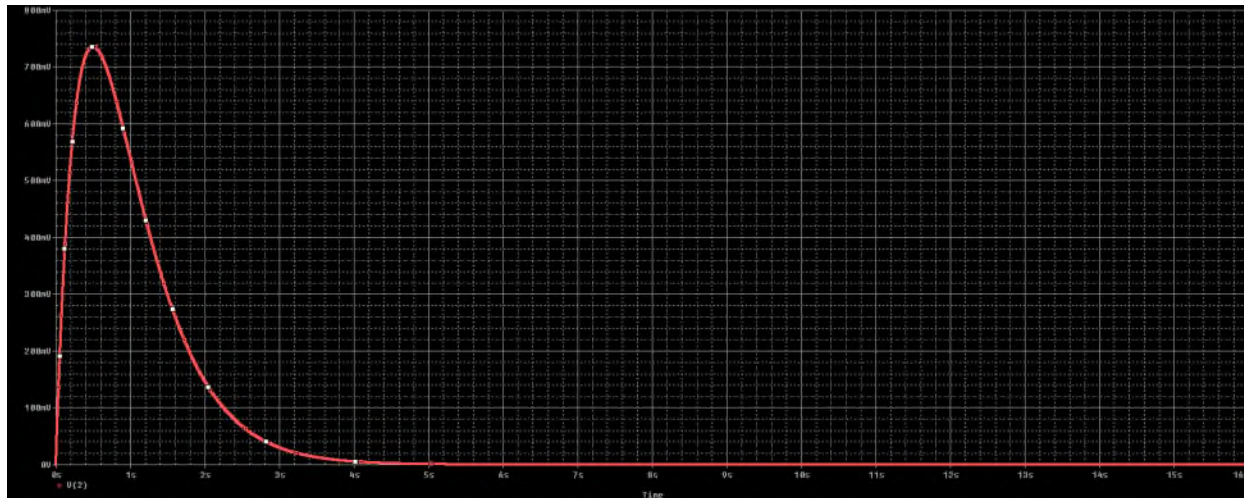
To see what this looks like, we can simulate the circuit with PSpice as follows:

```

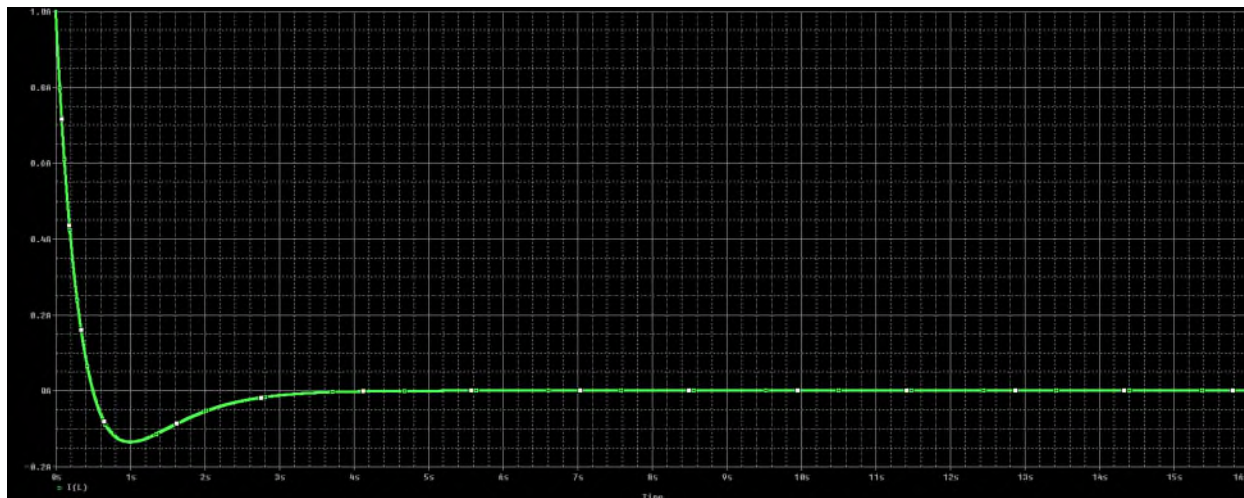
Example 3.2
L 0 1 1 IC=1
R 1 2 4
C 2 0 {1/4} IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END

```

The capacitor voltage is:



and the inductor current is:



Case 4

If $\zeta = 0$, then $\zeta^2 - 1 = -1$, and there will be two conjugate imaginary roots, $r_{1,2} = \pm j\omega_n$. In this case, the system is said to be *undamped*.

As there are *two* distinct roots to the characteristic equation, $y(t)$ has *two* exponential components

$$y(t) = \beta_1 e^{j\omega_n t} + \beta_2 e^{-j\omega_n t} \quad (1.111)$$

Here again, as in Case 2, it is usually preferred to use Euler's identity to express $y(t)$ in the alternate form

$$\begin{aligned} y(t) &= \beta_1 (\cos \omega_n t + j \sin \omega_n t) + \beta_2 (\cos \omega_n t - j \sin \omega_n t) \\ &= (\beta_1 + \beta_2) \cos \omega_n t + j(\beta_1 - \beta_2) \sin \omega_n t \\ &= B_1 \cos \omega_n t + B_2 \sin \omega_n t \end{aligned} \quad (1.112)$$

where $B_1 = \beta_1 + \beta_2$ and $B_2 = j(\beta_1 - \beta_2)$.

To determine the values of B_1 and B_2 note that

$$\dot{y}(t) = -\omega_n B_1 \sin \omega_n t + \omega_n B_2 \cos \omega_n t \quad (1.113)$$

Evaluating equations (1.112) and (1.113) at $t = 0$, we have

$$B_1 = y(0) \quad (1.114)$$

and

$$\omega_n B_2 = \dot{y}(0) \quad (1.115)$$

so that

$$B_2 = \frac{\dot{y}(0)}{\omega_n} \quad (1.116)$$

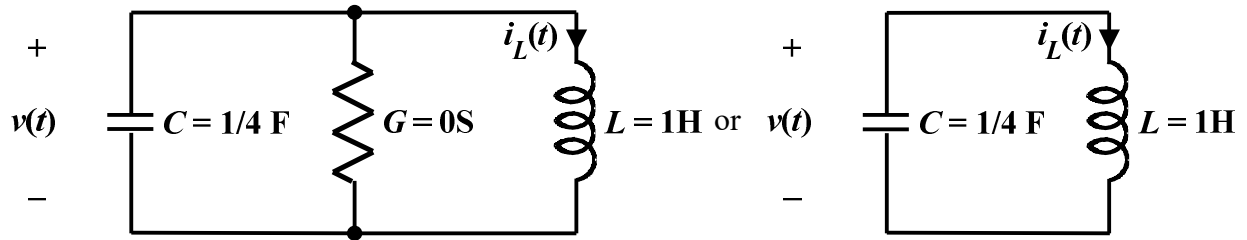
Alternately, note that

$$B_1 \cos \omega_n t + B_2 \sin \omega_n t = B_3 \cos(\omega_n t - \phi) \quad (1.117)$$

where $B_3 = \sqrt{B_1^2 + B_2^2}$ and $\phi = \tan^{-1}\left(\frac{B_2}{B_1}\right)$, so that $y(t)$ can be written in a slightly more compact form as

$$y(t) = B_3 \cos(\omega_n t - \phi) \quad (1.118)$$

Example 4.1



As shown by equation (1.7), this parallel RLC circuit can be described by the equation

$$\frac{d^2 i_L}{dt^2} + 4i_L = 0 \quad (1.119)$$

Hence, the characteristic equation is

$$r^2 + 4 = 0 \quad (1.120)$$

and

$$\omega_n = 2 \quad (1.121)$$

$$\zeta = 0 \quad (1.122)$$

This is an undamped system, with

$$r_1 = j2 \quad (1.123)$$

$$r_2 = -j2 \quad (1.124)$$

Suppose now that $i_L(0) = 0$ and $v(0) = 1$. Then, $v(t) = L \frac{di_L(t)}{dt}$, when evaluated at $t = 0$,

yields $\left. \frac{di_L}{dt} \right|_{t=0} = \frac{1}{L} v(0) = 1$, and

$$B_1 = 0 \quad (1.125)$$

$$B_2 = \frac{1}{2} \quad (1.126)$$

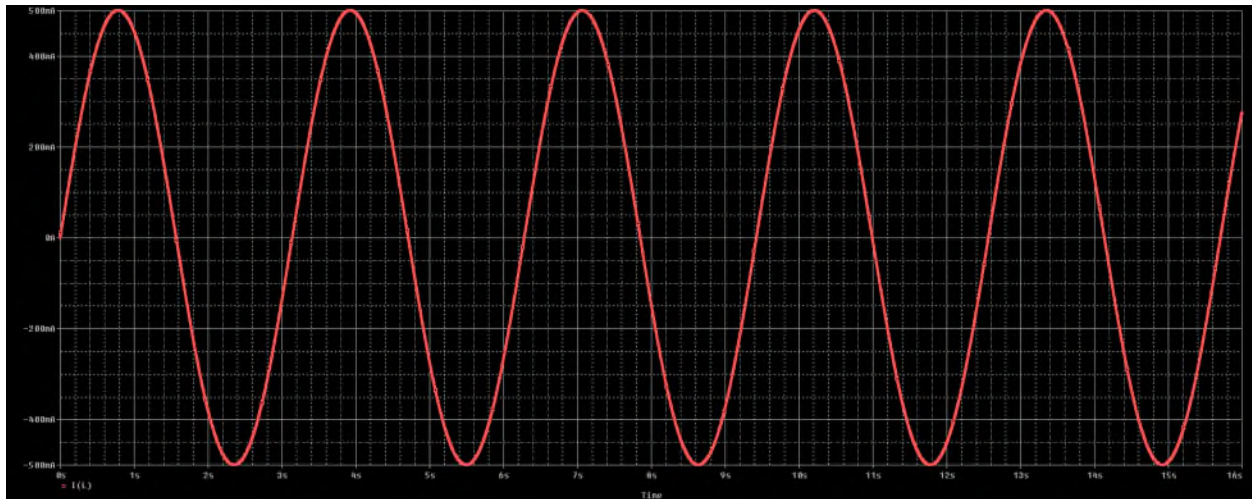
Hence,

$$i_L(t) = \frac{1}{2} \sin 2t \text{ A for } t > 0 \quad (1.127)$$

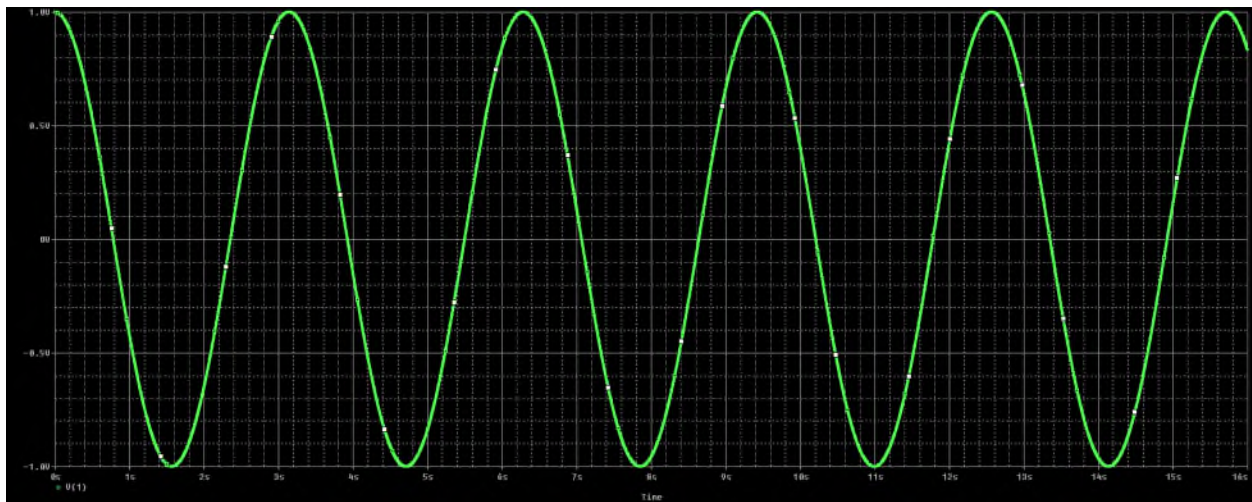
To see what this looks like, we can simulate the circuit with PSpice as follows:

```
Example 4.1
C 1 0 {1/4} IC=1
L 1 0 1 IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END
```

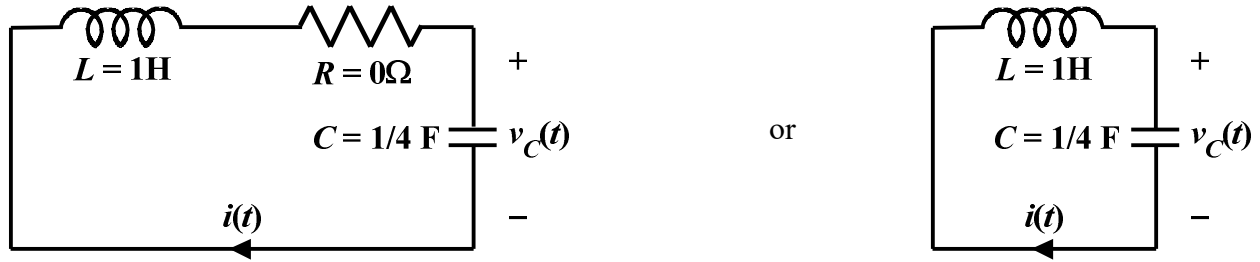
The inductor current is:



and the capacitor voltage is:



Example 4.2



As shown by equation (1.16), this series RLC circuit can be described by the equation

$$\frac{d^2 v_C}{dt^2} + 4v_C = 0 \quad (1.128)$$

Hence, the characteristic equation is

$$r^2 + 4 = 0 \quad (1.129)$$

and

$$\omega_n = 2 \quad (1.130)$$

$$\zeta = 0 \quad (1.131)$$

This is an undamped system, with

$$r_1 = j2 \quad (1.132)$$

$$r_2 = -j2 \quad (1.133)$$

Suppose now that $v_C(0) = 0$ and $i(0) = 1$. Then, $i(t) = C \frac{dv_C(t)}{dt}$, when evaluated at $t = 0$,

yields $\left. \frac{dv_C(t)}{dt} \right|_{t=0} = \frac{1}{C} i(0) = 4$, and

$$B_1 = 0 \quad (1.134)$$

$$B_2 = \frac{4}{2} = 2 \quad (1.135)$$

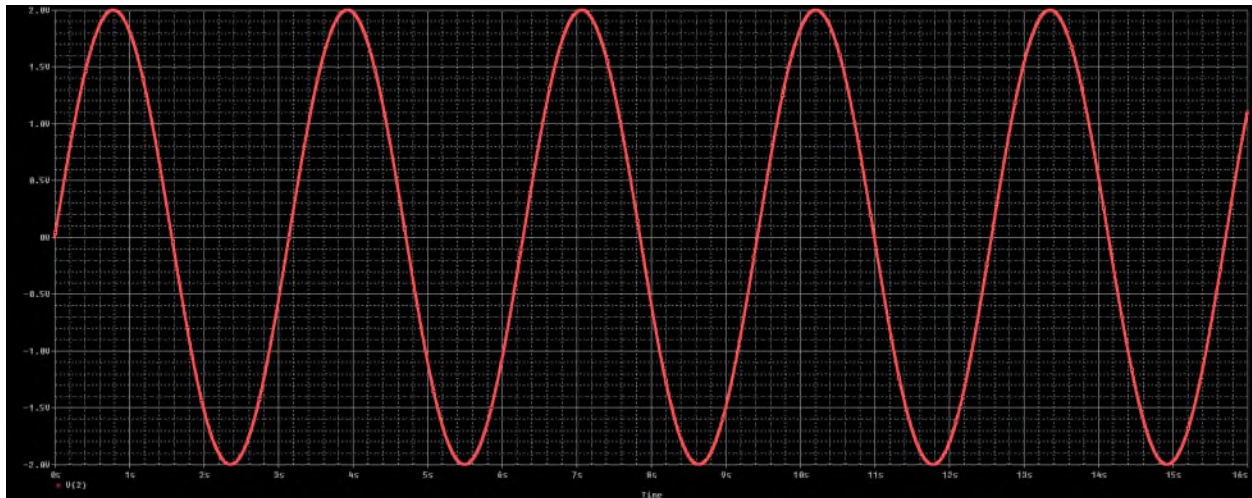
Hence,

$$v_C(t) = 2 \sin 2t \quad \text{V } t > 0 \quad (1.136)$$

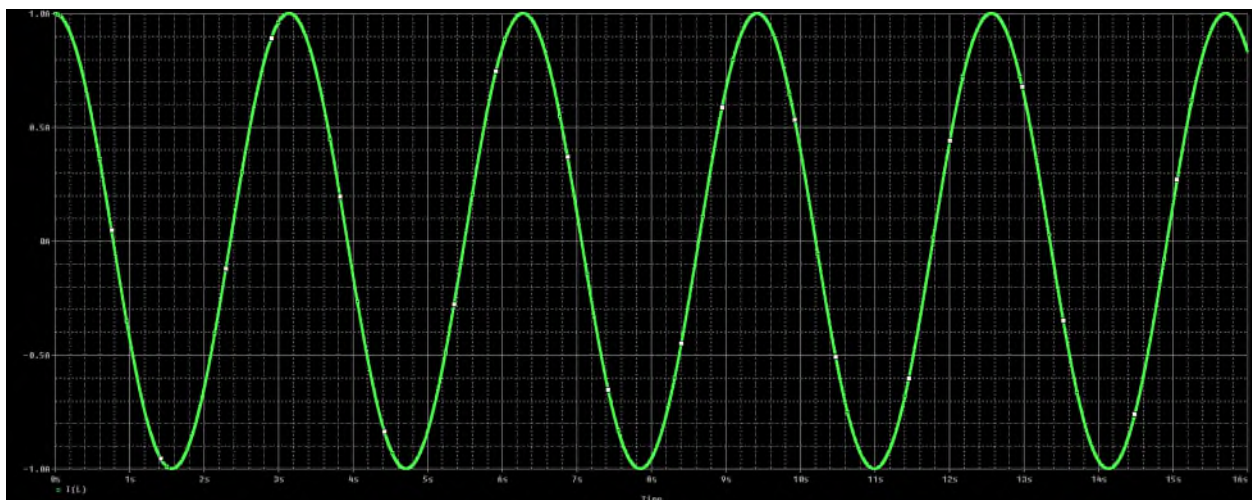
To see what this looks like, we can simulate the circuit with PSpice as follows:

```
Example 4.2
L 0 2 1 IC=1
C 2 0 {1/4} IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END
```

The capacitor voltage is:



and the inductor current is:

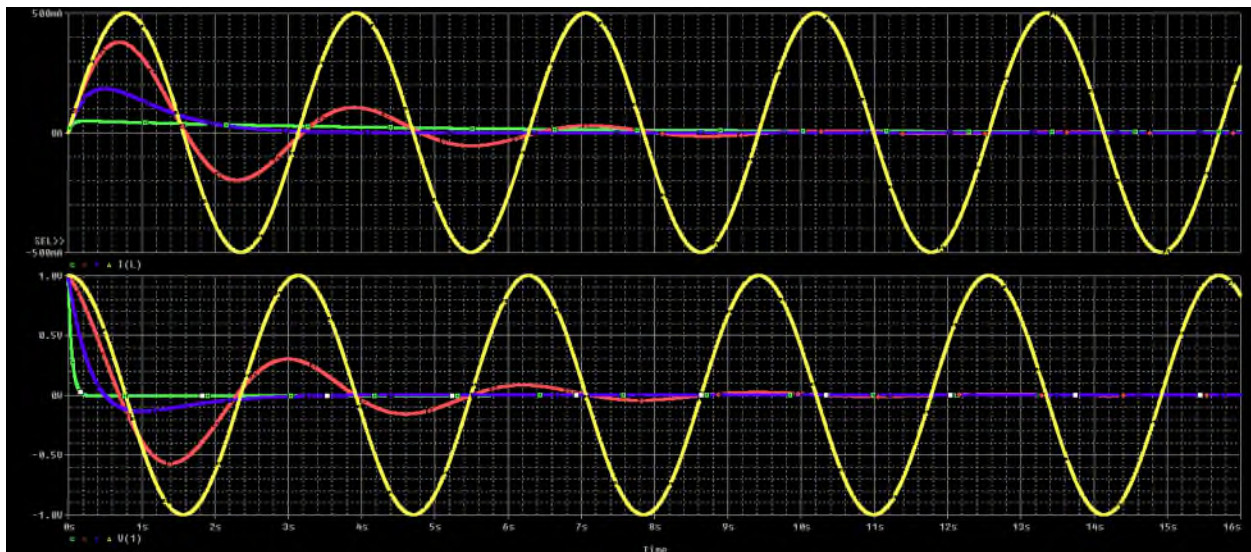


A comparison of the responses of the four parallel circuit examples (1.1, 2.1, 3.1 and 4.1) is shown below:

```

Example 1.1
C 1 0 {1/4} IC=1
G 1 0 1 0 5
L 1 0 1 IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END
Example 2.1
C 1 0 {1/4} IC=1
G 1 0 1 0 {1/5}
L 1 0 1 IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END
Example 3.1
C 1 0 {1/4} IC=1
G 1 0 1 0 1
L 1 0 1 IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END
Example 4.1
C 1 0 {1/4} IC=1
L 1 0 1 IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END

```

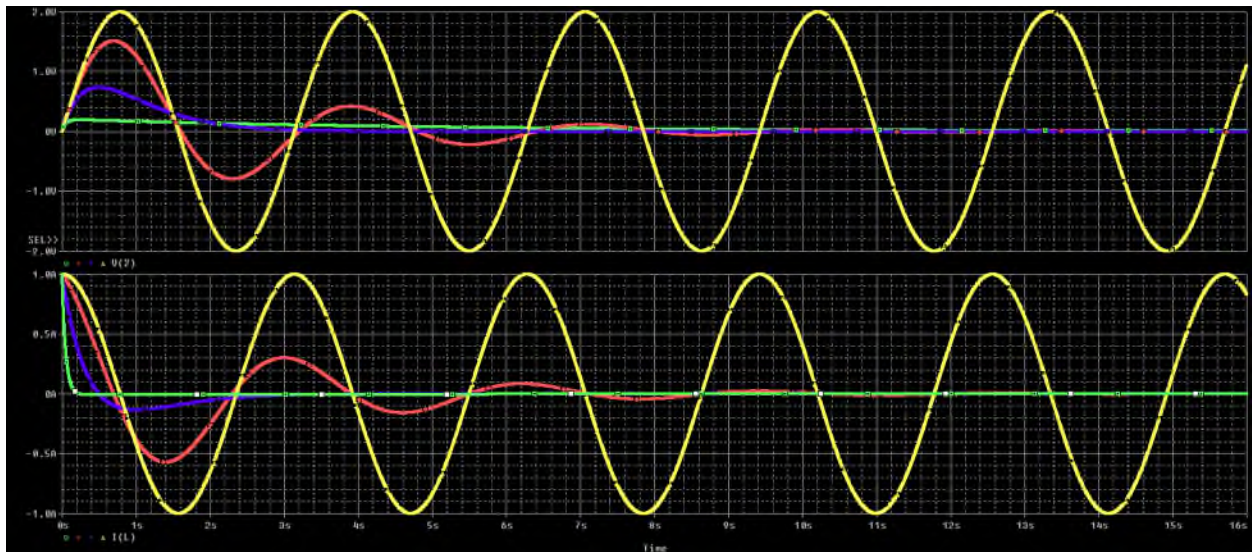


A comparison of the responses of the four series circuit examples (1.2, 2.2, 3.2 and 4.2) is shown below:

```

Example 1.2
L 0 1 1 IC=1
R 1 2 20
C 2 0 {1/4} IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END
Example 2.2
L 0 1 1 IC=1
R 1 2 {4/5}
C 2 0 {1/4} IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END
Example 3.2
L 0 1 1 IC=1
R 1 2 4
C 2 0 {1/4} IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END
Example 4.2
L 0 2 1 IC=1
C 2 0 {1/4} IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END

```



Systems with a Constant Input

Next consider systems with constant input, $z(t) = K$. In the case of electrical circuits, this means DC sources are applied. Equation (1.1) becomes:

$$\frac{d^2 y(t)}{dt^2} + 2\zeta\omega_n \frac{dy(t)}{dt} + \omega_n^2 y(t) = A\omega_n^2 K \quad (1.137)$$

If we assume that the natural response of the system is exponential, then $y(t) = \beta e^{rt} + \lambda$, and

$$r^2 \beta e^{rt} + 2\zeta\omega_n r \beta e^{rt} + \omega_n^2 (\beta e^{rt} + \lambda) = 0 \quad (1.138)$$

or

$$(r^2 + 2\zeta\omega_n r + \omega_n^2) \beta e^{rt} + \omega_n^2 \lambda = A\omega_n^2 K \quad (1.139)$$

which means that

$$r^2 + 2\zeta\omega_n r + \omega_n^2 = 0 \quad (1.140)$$

$$\omega_n^2 \lambda = A\omega_n^2 K \quad \Rightarrow \quad \lambda = AK \quad (1.141)$$

Equation (1.140) is called the characteristic equation of the system, and it has roots given by:

$$\begin{aligned} r_{1,2} &= \frac{-2\zeta\omega_n \pm \sqrt{(2\zeta\omega_n)^2 - 4\omega_n^2}}{2} \\ &= -\zeta\omega_n \pm \omega_n \sqrt{\zeta^2 - 1} \\ &= \left(-\zeta \pm \sqrt{\zeta^2 - 1}\right) \omega_n \end{aligned} \quad (1.142)$$

As in the unforced case, we will see that there are four distinctly different forms of the solution to equation (1.137), depending on the value of ζ with respect to the number 1.

Case 1

If $\zeta > 1$, then $\zeta^2 - 1 > 0$, and there will be two distinct negative real roots, $r_1 = (-\zeta + \sqrt{\zeta^2 - 1})\omega_n$ and $r_2 = (-\zeta - \sqrt{\zeta^2 - 1})\omega_n$. In this case, the system is said to be **overdamped**, and because there are **two** roots to the characteristic equation, $y(t)$ will have **two** exponential components:

$$y(t) = \beta_1 e^{r_1 t} + \beta_2 e^{r_2 t} + AK \quad (1.143)$$

To determine the values of β_1 and β_2 note that

$$\dot{y}(t) = \beta_1 r_1 e^{r_1 t} + \beta_2 r_2 e^{r_2 t} \quad (1.144)$$

Evaluating equations (1.143) and (1.144) at $t = 0$, we have

$$\beta_1 + \beta_2 + AK = y(0) \quad (1.145)$$

and

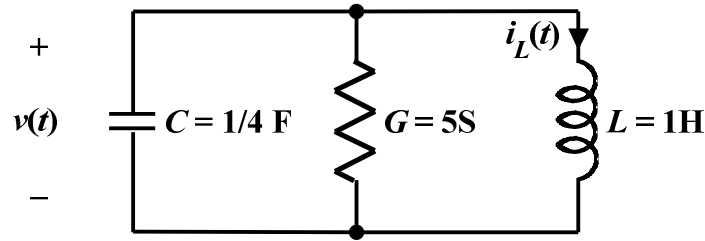
$$r_1 \beta_1 + r_2 \beta_2 = \dot{y}(0) \quad (1.146)$$

These two simultaneous equations can be used to evaluate β_1 and β_2 using Cramer's Rule as follows:

$$\beta_1 = \frac{\begin{vmatrix} y(0) - AK & 1 \\ \dot{y}(0) & r_2 \end{vmatrix}}{\begin{vmatrix} 1 & 1 \\ r_1 & r_2 \end{vmatrix}} = \frac{r_2 [y(0) - AK] - \dot{y}(0)}{r_2 - r_1} \quad (1.147)$$

$$\beta_2 = \frac{\begin{vmatrix} 1 & y(0) - AK \\ r_1 & \dot{y}(0) \end{vmatrix}}{\begin{vmatrix} 1 & 1 \\ r_1 & r_2 \end{vmatrix}} = \frac{\dot{y}(0) - r_1 [y(0) - AK]}{r_2 - r_1} \quad (1.148)$$

Example 1.1



As shown by equation (1.7), this parallel RLC circuit can be described by the equation

$$\frac{d^2 i_L}{dt^2} + 20 \frac{di_L}{dt} + 4i_L = 0 \quad (1.149)$$

Hence, the characteristic equation is

$$r^2 + 20r + 4 = 0 \quad (1.150)$$

and

$$\omega_n = \frac{1}{\sqrt{1(1/4)}} = 2 \quad (1.151)$$

$$\zeta = \frac{5}{2} \sqrt{\frac{1}{(1/4)}} = 5 \quad (1.152)$$

This is an overdamped system, with

$$r_1 = (-5 + \sqrt{25-1})2 \approx -0.202 \quad (1.153)$$

$$r_2 = (-5 - \sqrt{25-1})2 \approx -19.798 \quad (1.154)$$

Suppose now that $i_L(0) = 0$ and $v(0) = 1$. Then, $v(t) = L \frac{di_L(t)}{dt}$, when evaluated at $t = 0$,

yields $\left. \frac{di_L}{dt} \right|_{t=0} = \frac{1}{L} v(0) = 1$, and

$$\beta_1 \approx \frac{(-19.798)(0) - 1}{-19.798 - (-0.202)} \approx \frac{-1}{-19.596} \approx 0.051 \quad (1.155)$$

$$\beta_2 \approx \frac{1 - (-0.202)(0)}{-19.798 - (-0.202)} \approx \frac{1}{-19.596} \approx -0.051 \quad (1.156)$$

Hence,

$$i_L(t) \approx 0.051e^{-0.202t} - 0.051e^{-19.798t} \text{ A for } t > 0 \quad (1.157)$$

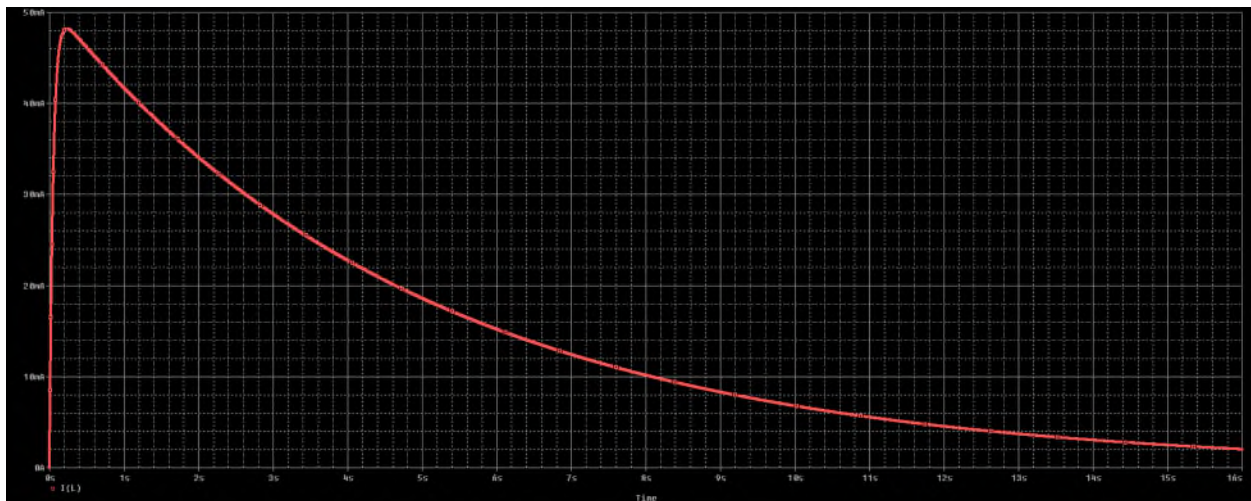
To see what this looks like, we can simulate the circuit with PSpice as follows:

```

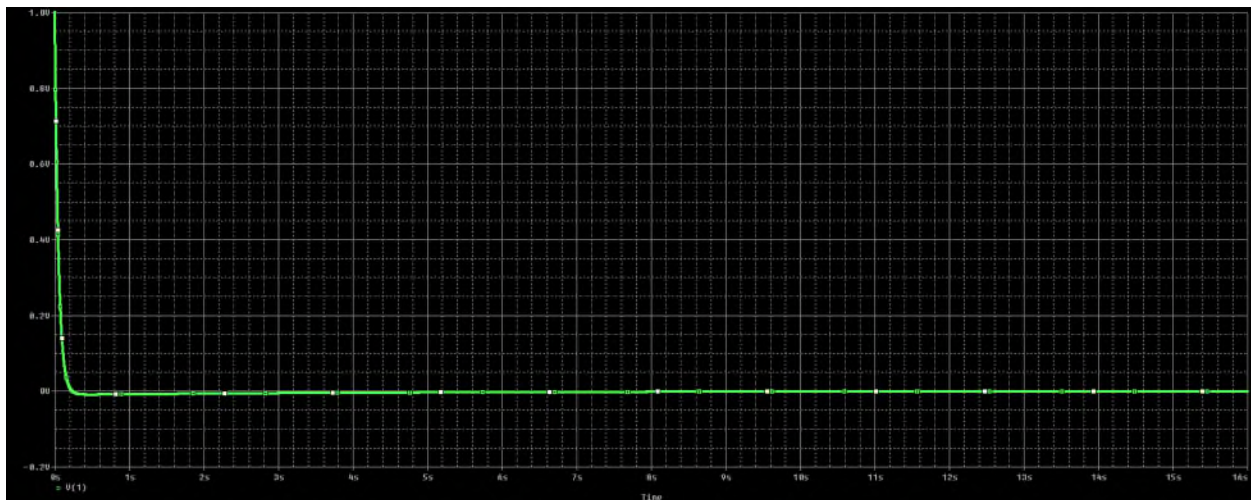
Example 1.1
C      1      0      {1/4}      IC=1
G      1      0      1      0      5
L      1      0      1      IC=0
.TRAN  1      16      0      1m      UIC
.PROBE
.END

```

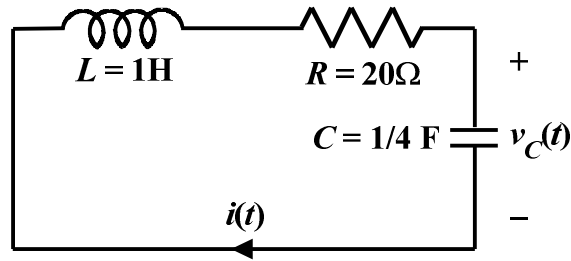
The inductor current is:



and the capacitor voltage is:



Example 1.2



As shown by equation (1.16), this series RLC circuit can be described by the equation

$$\frac{d^2 v_C}{dt^2} + 20 \frac{dv_C}{dt} + 4v_C = 0 \quad (1.158)$$

Hence, the characteristic equation is

$$r^2 + 20r + 4 = 0 \quad (1.159)$$

and

$$\omega_n = \frac{1}{\sqrt{1(1/4)}} = 2 \quad (1.160)$$

$$\zeta = \frac{20}{2} \sqrt{\frac{(1/4)}{1}} = 5 \quad (1.161)$$

This is an overdamped system, with

$$r_1 = (-5 + \sqrt{25-1})2 \approx -0.202 \quad (1.162)$$

$$r_2 = (-5 - \sqrt{25-1})2 \approx -19.798 \quad (1.163)$$

Suppose now that $v_C(0) = 0$ and $i(0) = 1$. Then, $i(t) = C \frac{dv_C(t)}{dt}$, when evaluated at $t = 0$,

yields $\left. \frac{dv_C(t)}{dt} \right|_{t=0} = \frac{1}{C} i(0) = 4$, and

$$\beta_1 \approx \frac{(-19.798)(0) - 4}{-19.798 - (-0.202)} \approx 0.204 \quad (1.164)$$

$$\beta_2 \approx \frac{4 - (-0.202)(0)}{-19.798 - (-0.202)} \approx -0.204 \quad (1.165)$$

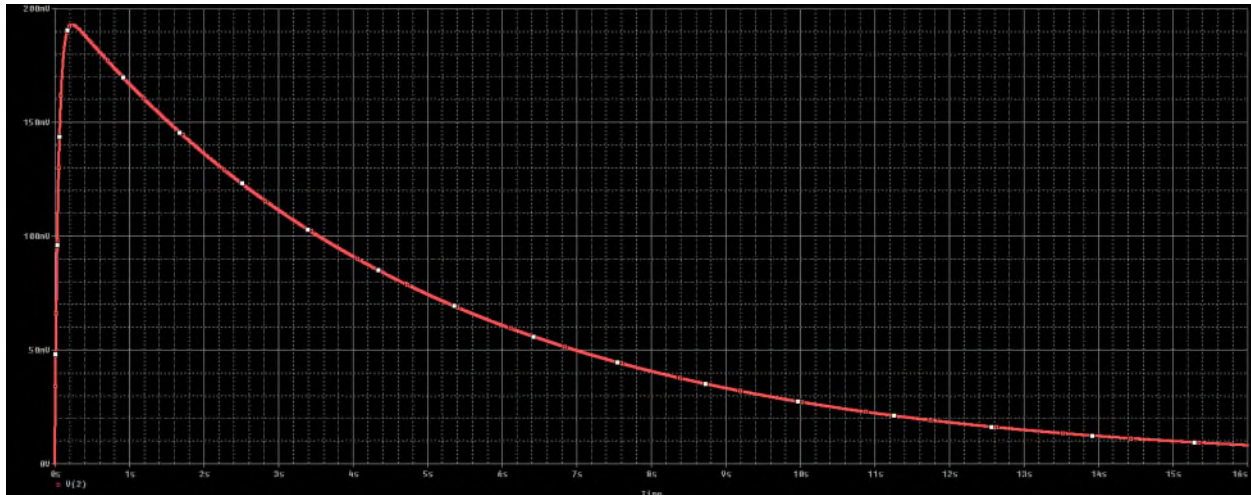
Hence,

$$v_C(t) \approx 0.204e^{-0.202t} - 0.204e^{-19.798t} \text{ V for } t > 0 \quad (1.166)$$

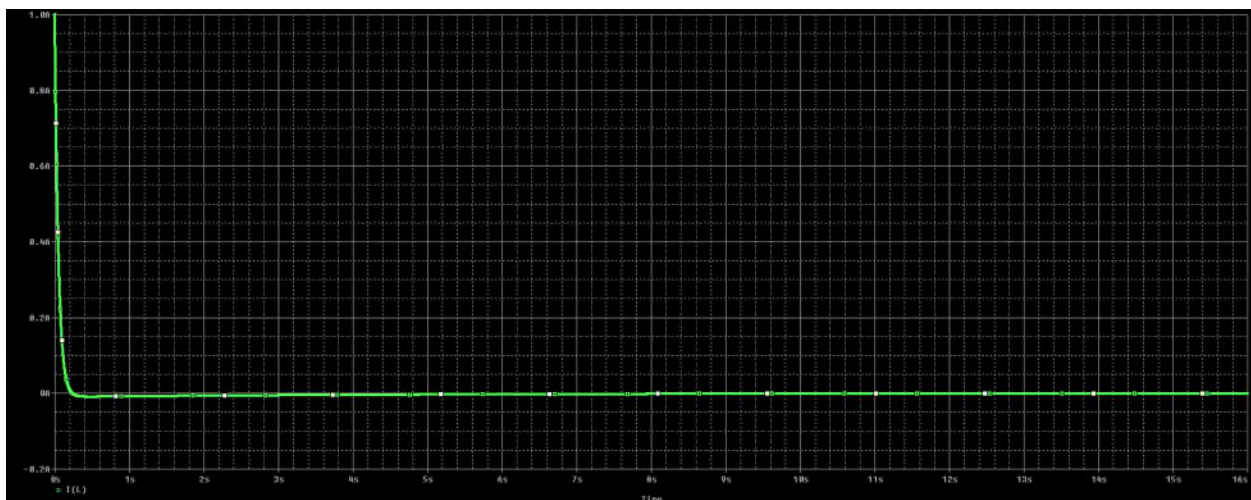
To see what this looks like, we can simulate the circuit with PSpice as follows:

```
Example 1.2
L      0      1      1      IC=1
R      1      2      20
C      2      0      {1/4}    IC=0
.TRAN  1      16     0      1m    UIC
.PROBE
.END
```

The capacitor voltage is:



and the inductor current is:



Case 2

If $0 < \zeta < 1$, then $\zeta^2 - 1 < 0$, and there will be two complex conjugate roots,

$r_1 = \left(-\zeta + j\sqrt{1-\zeta^2}\right)\omega_n = -\zeta\omega_n + j\omega_d$ and $r_2 = \left(-\zeta - j\sqrt{1-\zeta^2}\right)\omega_n = -\zeta\omega_n - j\omega_d$. In this case,

the system is said to be **underdamped**., and the quantity $\omega_d = \omega_n\sqrt{1-\zeta^2}$ is called the *damped* or *ringing* frequency.

As in Case 1, because there are **two** distinct roots to the characteristic equation, $y(t)$ has **two** exponential components:

$$\begin{aligned} y(t) &= \beta_1 e^{(-\zeta\omega_n + j\omega_d)t} + \beta_2 e^{(-\zeta\omega_n - j\omega_d)t} + AK \\ &= e^{-\zeta\omega_n t} (\beta_1 e^{j\omega_d t} + \beta_2 e^{-j\omega_d t}) + AK \end{aligned} \quad (1.167)$$

However, it is usually preferred to use Euler's identity

$$e^{\pm j\theta} = \cos \theta \pm j \sin \theta \quad (1.168)$$

to express $y(t)$ in the alternate form

$$\begin{aligned} y(t) &= e^{-\zeta\omega_n t} [\beta_1 (\cos \omega_d t + j \sin \omega_d t) + \beta_2 (\cos \omega_d t - j \sin \omega_d t)] + AK \\ &= e^{-\zeta\omega_n t} [(\beta_1 + \beta_2) \cos \omega_d t + j(\beta_1 - \beta_2) \sin \omega_d t] + AK \\ &= e^{-\zeta\omega_n t} [B_1 \cos \omega_d t + B_2 \sin \omega_d t] + AK \end{aligned} \quad (1.169)$$

where $B_1 = \beta_1 + \beta_2$ and $B_2 = j(\beta_1 - \beta_2)$.

To determine the values of B_1 and B_2 note that

$$\dot{y}(t) = -\zeta\omega_n e^{-\zeta\omega_n t} [B_1 \cos \omega_d t + B_2 \sin \omega_d t] + e^{-\zeta\omega_n t} [-B_1\omega_d \sin \omega_d t + B_2\omega_d \cos \omega_d t] \quad (1.170)$$

Evaluating equations (1.169) and (1.170) at $t = 0$, we have

$$B_1 + AK = y(0) \quad (1.171)$$

and

$$-\zeta\omega_n B_1 + B_2\omega_d = \dot{y}(0) \quad (1.172)$$

Thus,

$$B_1 = y(0) - AK \quad (1.173)$$

and

$$B_2 = \frac{\dot{y}(0) + \zeta\omega_n B_1}{\omega_d} = \frac{\dot{y}(0) + \zeta\omega_n [y(0) - AK]}{\omega_d} \quad (1.174)$$

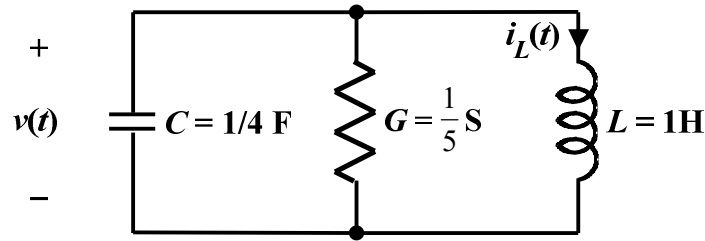
Alternately, note that

$$B_1 \cos \omega_d t + B_2 \sin \omega_d t = B_3 \cos(\omega_d t - \phi) \quad (1.175)$$

where $B_3 = \sqrt{B_1^2 + B_2^2}$ and $\phi = \tan^{-1}\left(\frac{B_2}{B_1}\right)$, so that $y(t)$ can be written in a slightly more compact form as

$$y(t) = B_3 e^{-\zeta\omega_n t} \cos(\omega_d t - \phi) + AK \quad (1.176)$$

Example 2.1



As shown by equation (1.7), this parallel RLC circuit can be described by the equation

$$\frac{d^2 i_L}{dt^2} + \frac{4}{5} \frac{di_L}{dt} + 4i_L = 0 \quad (1.177)$$

Hence, the characteristic equation is

$$r^2 + \frac{4}{5}r + 4 = 0 \quad (1.178)$$

and

$$\omega_n = 2 \quad (1.179)$$

$$\zeta = \frac{1}{10} \sqrt{4} = 0.2 \quad (1.180)$$

$$\omega_d = 2\sqrt{1 - (0.2)^2} \approx 1.960 \quad (1.181)$$

This is an underdamped system, with

$$r_{1,2} \approx -0.400 + j1.960 \quad (1.182)$$

$$r_{1,2} \approx -0.400 - j1.960 \quad (1.183)$$

Suppose now that $i_L(0) = 0$ and $v(0) = 1$. Then, $v(t) = L \frac{di_L(t)}{dt}$, when evaluated at $t = 0$,

yields $\left. \frac{di_L(t)}{dt} \right|_{t=0} = \frac{1}{L} v(0) = 1$, and

$$B_1 = 0 \quad (1.184)$$

$$B_2 \approx \frac{1 + (0.2)(2)(0)}{1.960} \approx 0.510 \quad (1.185)$$

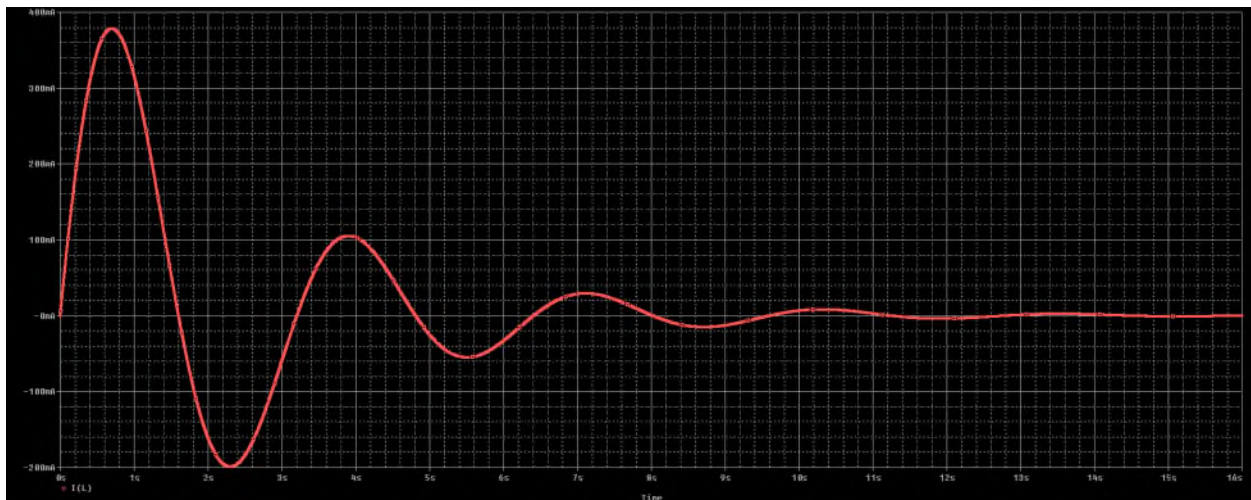
Hence,

$$i_L(t) \approx 0.510e^{-0.4t} \sin(1.960t) \text{ A for } t > 0 \quad (1.186)$$

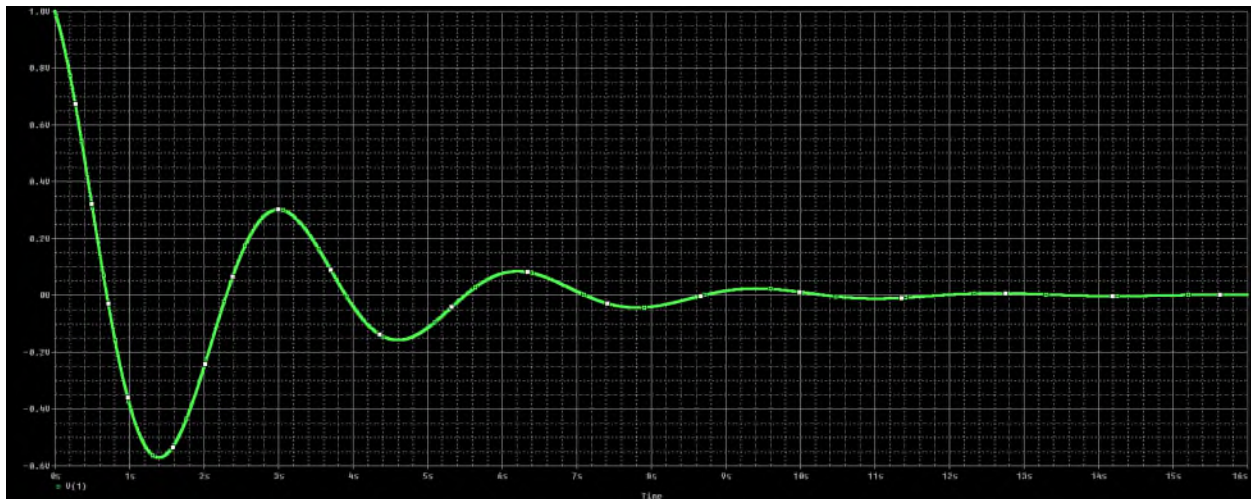
To see what this looks like, we can simulate the circuit with PSpice as follows:

```
Example 2.1
C 1 0 {1/4} IC=1
G 1 0 1 0 {1/5}
L 1 0 1 IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END
```

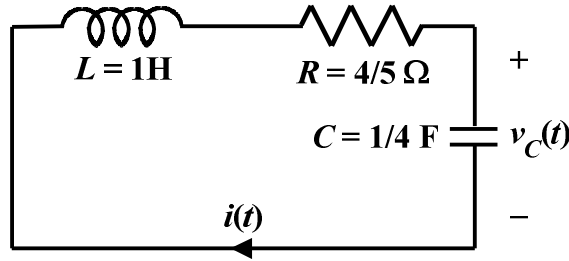
The inductor current is:



and the capacitor voltage is:



Example 2.2



As shown by equation (1.16), this series RLC circuit can be described by the equation

$$\frac{d^2 v_C}{dt^2} + \frac{4}{5} \frac{dv_C}{dt} + 4v_C = 0 \quad (1.187)$$

Hence, the characteristic equation is

$$r^2 + \frac{4}{5}r + 4 = 0 \quad (1.188)$$

and

$$\omega_n = 2 \quad (1.189)$$

$$\zeta = \frac{2}{5} \sqrt{\frac{1}{4}} = 0.2 \quad (1.190)$$

$$\omega_d = 2\sqrt{1 - (0.2)^2} \approx 1.960 \quad (1.191)$$

This is an underdamped system, with

$$r_1 = -0.400 + j1.960 \quad (1.192)$$

$$r_2 = -0.400 - j1.960 \quad (1.193)$$

Suppose now that $v_C(0) = 0$ and $i(0) = 1$. Then, $i(t) = C \frac{dv_C(t)}{dt}$, when evaluated at $t = 0$,

yields $\left. \frac{dv_C(t)}{dt} \right|_{t=0} = \frac{1}{C} i(0) = 4$, and

$$B_1 = 0 \quad (1.194)$$

$$B_2 \approx \frac{4 + (0.2)(2)(0)}{1.960} \approx 2.041 \quad (1.195)$$

Hence,

$$v_c(t) = 2.041e^{-0.4t} \sin(1.960t) \quad \forall t > 0 \quad (1.196)$$

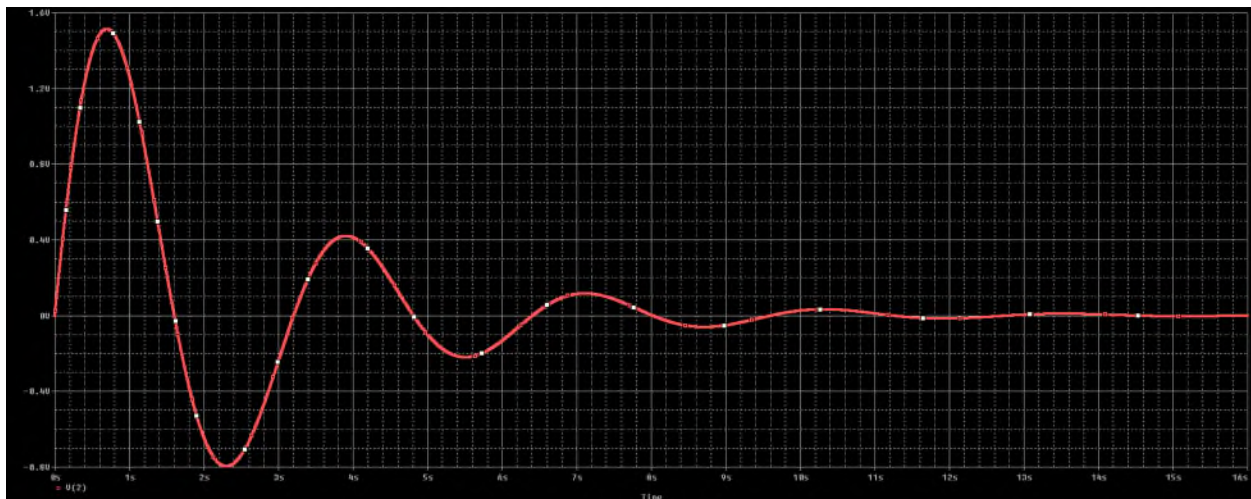
To see what this looks like, we can simulate the circuit with PSpice as follows:

```

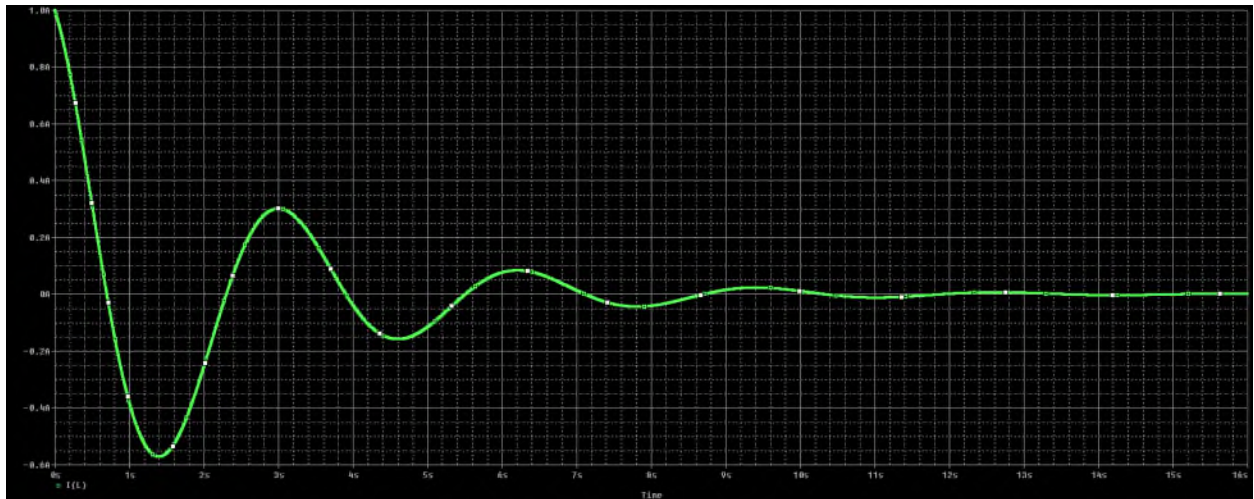
Example 2.2
L      0      1      1      IC=1
R      1      2      {4/5}
C      2      0      {1/4}      IC=0
.TRAN  1      16      0      1m      UIC
.PROBE
.END

```

The capacitor voltage is:



and the inductor current is:



Case 3

If $\zeta = 1$, then $\zeta^2 - 1 = 0$, and there will be two identical negative real roots, $r_1 = r_2 = -\omega_n$. In this case, the system is said to be *critically damped*. This case can be considered to be the “borderline” between overdamped and underdamped systems.

The general form of the solution is

$$y(t) = (\beta_1 + \beta_2 t)e^{-\omega_n t} + AK \quad (1.197)$$

To determine the values of β_1 and β_2 note that

$$\dot{y}(t) = \beta_2 e^{-\omega_n t} - \omega_n (\beta_1 + \beta_2 t)e^{-\omega_n t} \quad (1.198)$$

Evaluating equations (1.197) and (1.198) at $t = 0$, we have

$$\beta_1 + AK = y(0) \quad (1.199)$$

and

$$\beta_2 - \omega_n \beta_1 = \dot{y}(0) \quad (1.200)$$

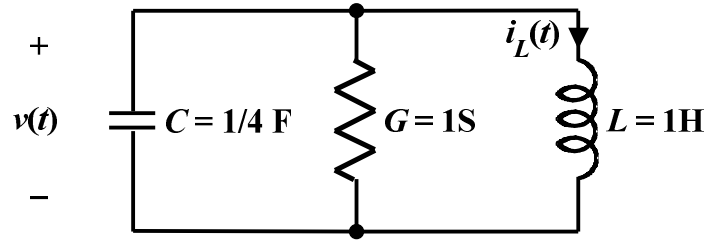
Thus,

$$\beta_1 = y(0) - AK \quad (1.201)$$

and

$$\beta_2 = \dot{y}(0) + \omega_n \beta_1 = \dot{y}(0) + \omega_n [y(0) - AK] \quad (1.202)$$

Example 3.1



As shown by equation (1.7), this parallel circuit can be described by the equation

$$\frac{d^2 i_L}{dt^2} + 4 \frac{di_L}{dt} + 4i_L = 0 \quad (1.203)$$

Hence, the characteristic equation is

$$r^2 + 4r + 4 = 0 \quad (1.204)$$

and

$$\omega_n = 2 \quad (1.205)$$

$$\zeta = \frac{1}{2} \sqrt{4} = 1 \quad (1.206)$$

This is a critically damped system, with

$$r_1 = r_2 = -2 \quad (1.207)$$

Suppose now that $i_L(0) = 0$ and $v(0) = 1$. Then, $v(t) = L \frac{di_L(t)}{dt}$, when evaluated at $t = 0$,

yields $\left. \frac{di_L}{dt} \right|_{t=0} = \frac{1}{L} v(0) = 1$, and

$$\beta_1 = 0 \quad (1.208)$$

$$\beta_2 = 1 + (2)(0) = 1 \quad (1.209)$$

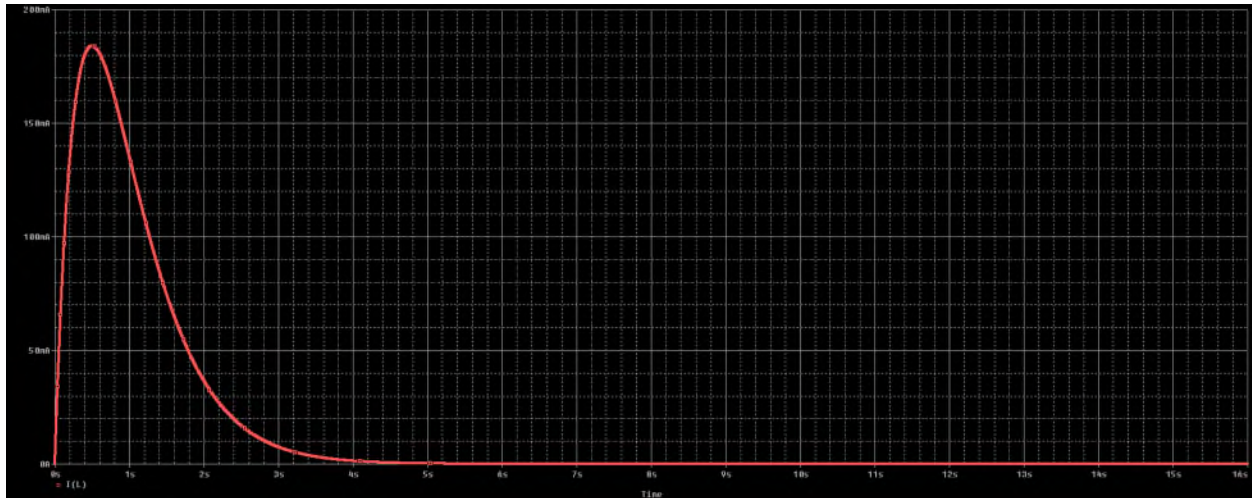
Hence,

$$i_L(t) = te^{-2t} \text{ A for } t > 0 \quad (1.210)$$

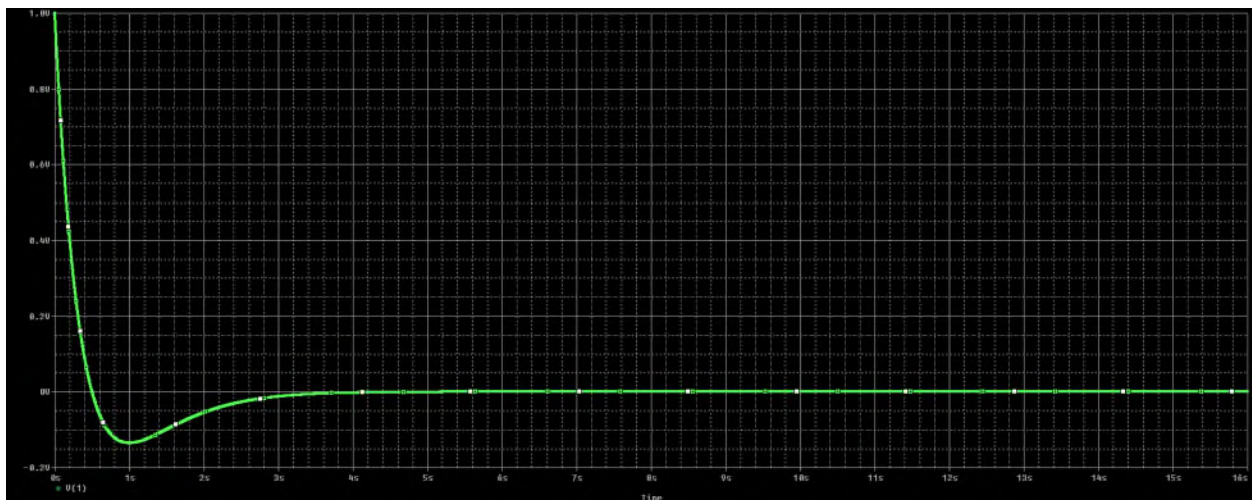
To see what this looks like, we can simulate the circuit with PSpice as follows:

```
Example 3.1
C 1 0 {1/4} IC=1
G 1 0 1 0 1
L 1 0 1 IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END
```

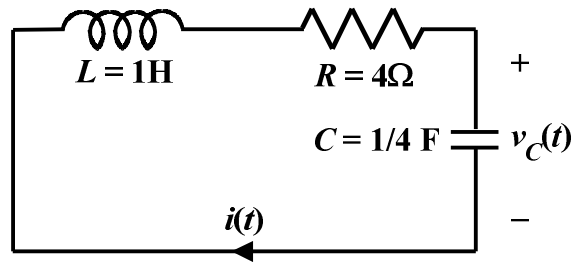
The inductor current is:



and the capacitor voltage is:



Example 3.2



As shown by equation (1.16), this series RLC circuit can be described by the equation

$$\frac{d^2v_C}{dt^2} + 4\frac{dv_C}{dt} + 4v_C = 0 \quad (1.211)$$

Hence, the characteristic equation is

$$r^2 + 4r + 4 = 0 \quad (1.212)$$

and

$$\omega_n = 2 \quad (1.213)$$

$$\zeta = \frac{4}{2}\sqrt{\frac{1}{4}} = 1 \quad (1.214)$$

This is a critically damped system, with

$$r_1 = r_2 = -2 \quad (1.215)$$

Suppose now that $v_C(0) = 0$ and $i(0) = 1$. Then, $i(t) = C\frac{dv_C(t)}{dt}$, when evaluated at $t = 0$,

yields $\left.\frac{dv_C(t)}{dt}\right|_{t=0} = \frac{1}{C}i(0) = 4$, and

$$\beta_1 = 0 \quad (1.216)$$

$$\beta_2 = 4 + (2)(0) = 4 \quad (1.217)$$

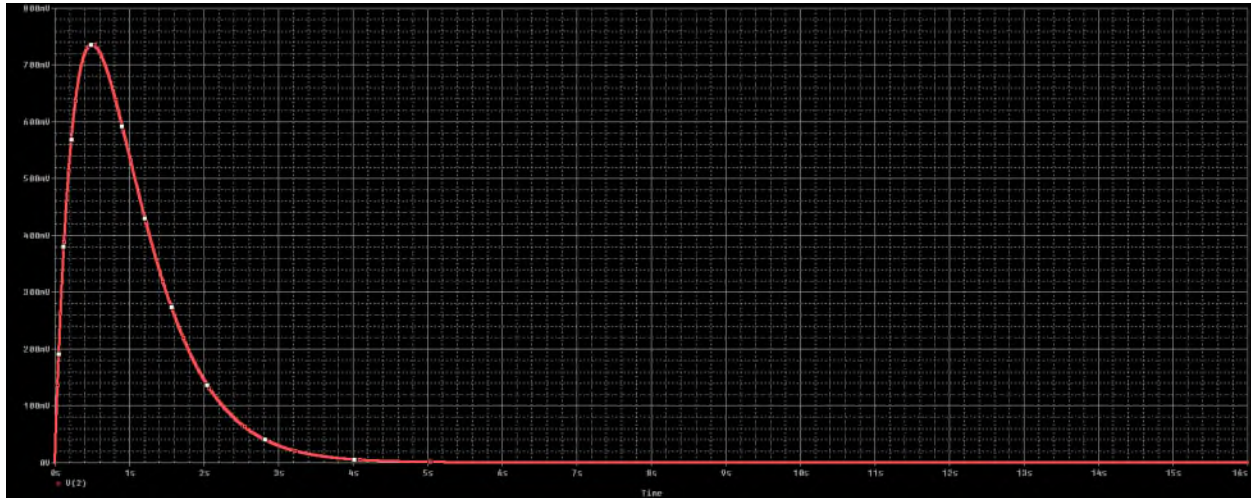
Hence,

$$v_C(t) = 4te^{-2t} \text{ V } t > 0 \quad (1.218)$$

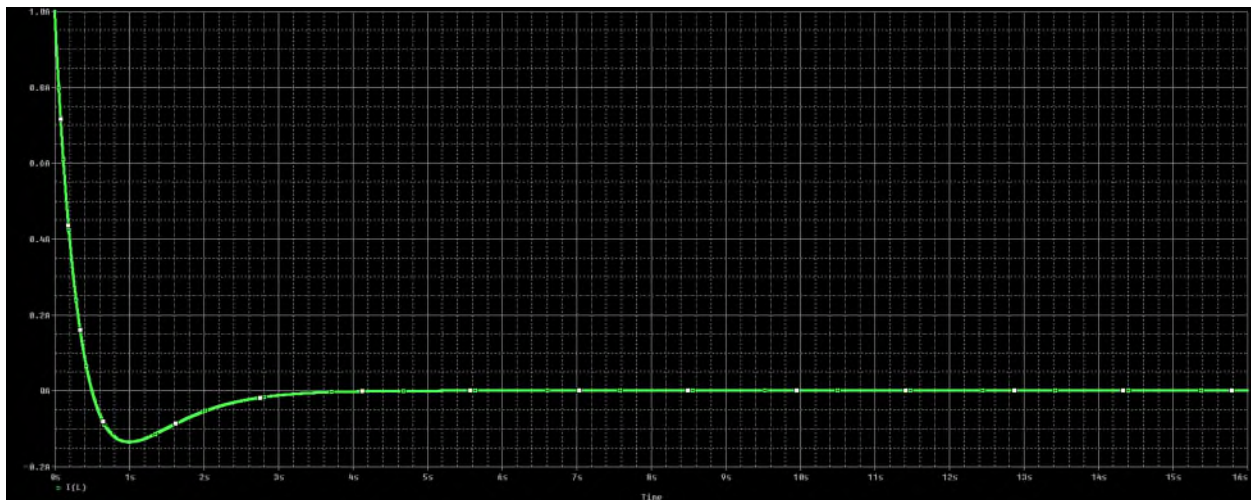
To see what this looks like, we can simulate the circuit with PSpice as follows:

```
Example 3.2
L 0 1 1 IC=1
R 1 2 4
C 2 0 {1/4} IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END
```

The capacitor voltage is:



and the inductor current is:



Case 4

If $\zeta = 0$, then $\zeta^2 - 1 = -1$, and there will be two conjugate imaginary roots, $r_{1,2} = \pm j\omega_n$. In this case, the system is said to be *undamped*.

As there are *two* distinct roots to the characteristic equation, $y(t)$ has *two* exponential components

$$y(t) = \beta_1 e^{j\omega_n t} + \beta_2 e^{-j\omega_n t} + AK \quad (1.219)$$

Here again, as in Case 2, it is usually preferred to use Euler's identity to express $y(t)$ in the alternate form

$$\begin{aligned} y(t) &= \beta_1 (\cos \omega_n t + j \sin \omega_n t) + \beta_2 (\cos \omega_n t - j \sin \omega_n t) + AK \\ &= (\beta_1 + \beta_2) \cos \omega_n t + j(\beta_1 - \beta_2) \sin \omega_n t + AK \\ &= B_1 \cos \omega_n t + B_2 \sin \omega_n t + AK \end{aligned} \quad (1.220)$$

where $B_1 = \beta_1 + \beta_2$ and $B_2 = j(\beta_1 - \beta_2)$.

To determine the values of B_1 and B_2 note that

$$\dot{y}(t) = -\omega_n B_1 \sin \omega_n t + \omega_n B_2 \cos \omega_n t \quad (1.221)$$

Evaluating equations (1.220) and (1.221) at $t = 0$, we have

$$B_1 + AK = y(0) \quad (1.222)$$

and

$$\omega_n B_2 = \dot{y}(0) \quad (1.223)$$

so that

$$B_1 = y(0) - AK \quad (1.224)$$

and

$$B_2 = \frac{\dot{y}(0)}{\omega_n} \quad (1.225)$$

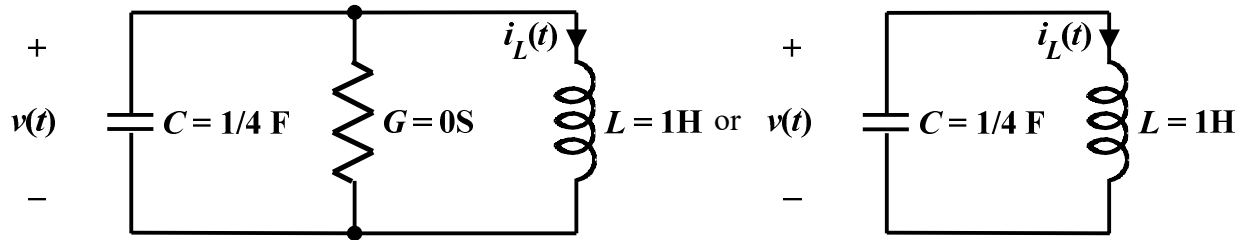
Alternately, note that

$$B_1 \cos \omega_n t + B_2 \sin \omega_n t = B_3 \cos(\omega_n t - \phi) \quad (1.226)$$

where $B_3 = \sqrt{B_1^2 + B_2^2}$ and $\phi = \tan^{-1}\left(\frac{B_2}{B_1}\right)$, so that $y(t)$ can be written in a slightly more compact form as

$$y(t) = B_3 \cos(\omega_n t - \phi) + AK \quad (1.227)$$

Example 4.1



As shown by equation (1.7), this parallel RLC circuit can be described by the equation

$$\frac{d^2 i_L}{dt^2} + 4i_L = 0 \quad (1.228)$$

Hence, the characteristic equation is

$$r^2 + 4 = 0 \quad (1.229)$$

and

$$\omega_n = 2 \quad (1.230)$$

$$\zeta = 0 \quad (1.231)$$

This is an undamped system, with

$$r_1 = j2 \quad (1.232)$$

$$r_2 = -j2 \quad (1.233)$$

Suppose now that $i_L(0) = 0$ and $v(0) = 1$. Then, $v(t) = L \frac{di_L(t)}{dt}$, when evaluated at $t = 0$,

yields $\left. \frac{di_L}{dt} \right|_{t=0} = \frac{1}{L} v(0) = 1$, and

$$B_1 = 0 \quad (1.234)$$

$$B_2 = \frac{1}{2} \quad (1.235)$$

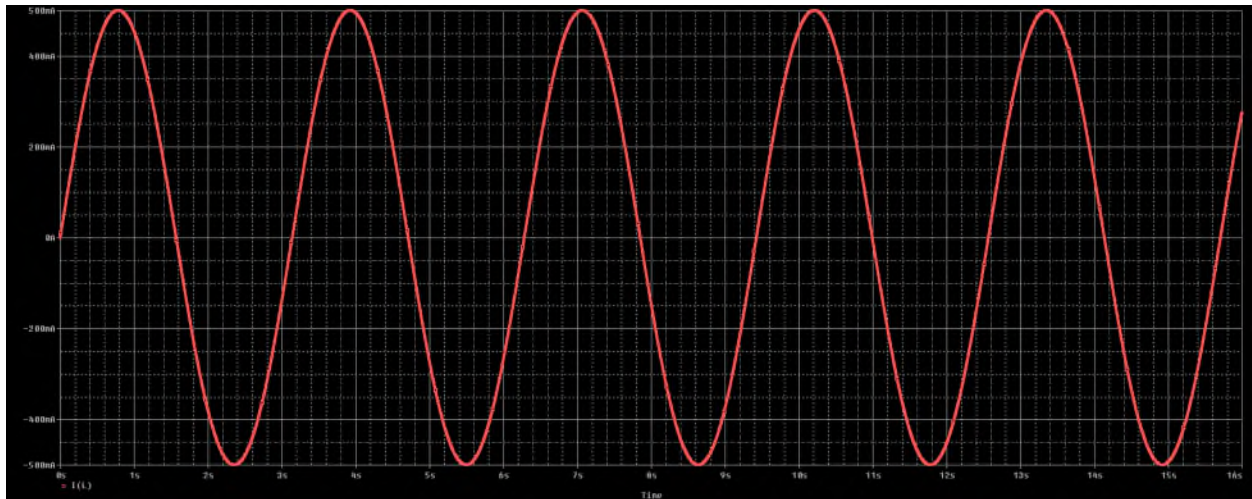
Hence,

$$i_L(t) = \frac{1}{2} \sin 2t \text{ A for } t > 0 \quad (1.236)$$

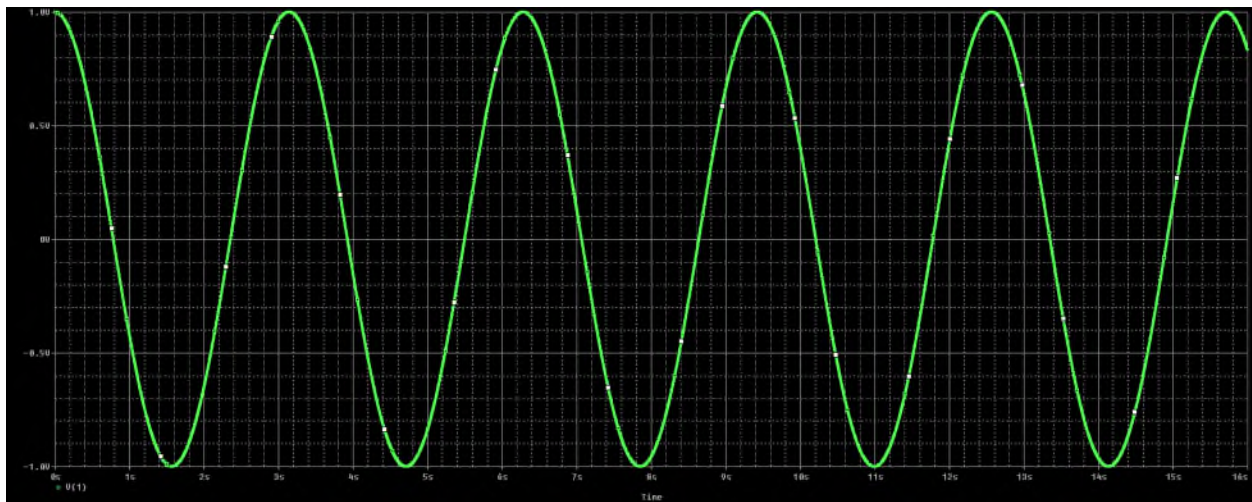
To see what this looks like, we can simulate the circuit with PSpice as follows:

```
Example 4.1
C 1 0 {1/4} IC=1
L 1 0 1 IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END
```

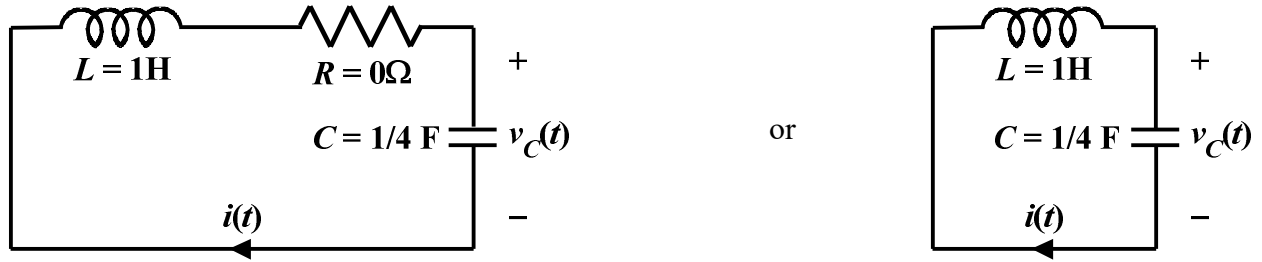
The inductor current is:



and the capacitor voltage is:



Example 4.2



As shown by equation (1.16), this series RLC circuit can be described by the equation

$$\frac{d^2 v_C}{dt^2} + 4v_C = 0 \quad (1.237)$$

Hence, the characteristic equation is

$$r^2 + 4 = 0 \quad (1.238)$$

and

$$\omega_n = 2 \quad (1.239)$$

$$\zeta = 0 \quad (1.240)$$

This is an undamped system, with

$$r_1 = j2 \quad (1.241)$$

$$r_2 = -j2 \quad (1.242)$$

Suppose now that $v_C(0) = 0$ and $i(0) = 1$. Then, $i(t) = C \frac{dv_C(t)}{dt}$, when evaluated at $t = 0$,

yields $\left. \frac{dv_C(t)}{dt} \right|_{t=0} = \frac{1}{C} i(0) = 4$, and

$$B_1 = 0 \quad (1.243)$$

$$B_2 = \frac{4}{2} = 2 \quad (1.244)$$

Hence,

$$v_C(t) = 2 \sin 2t \quad \text{V } t > 0 \quad (1.245)$$

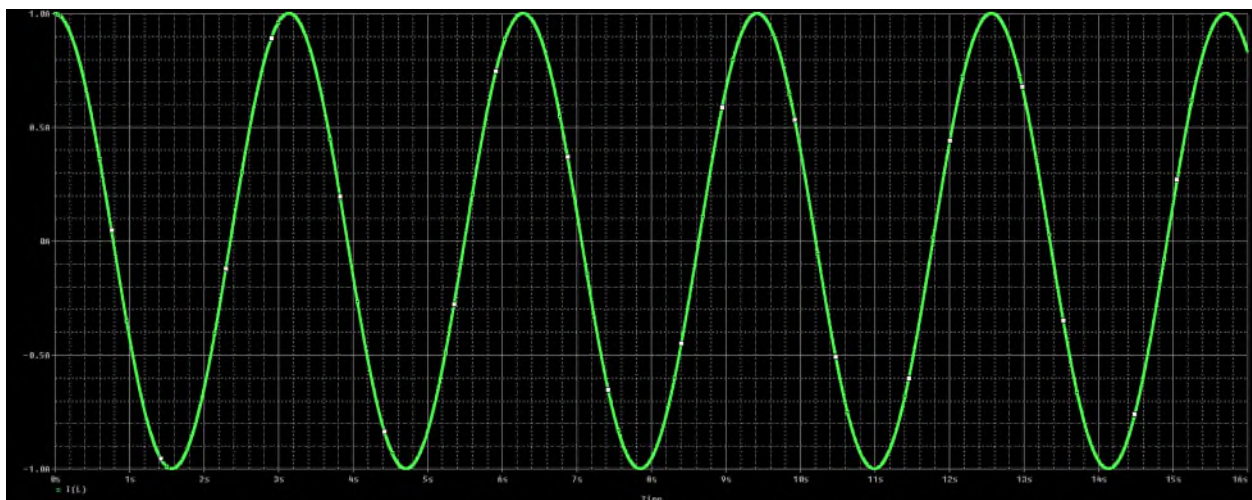
To see what this looks like, we can simulate the circuit with PSpice as follows:

```
Example 4.2
L 0 2 1 IC=1
C 2 0 {1/4} IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END
```

The capacitor voltage is:



and the inductor current is:

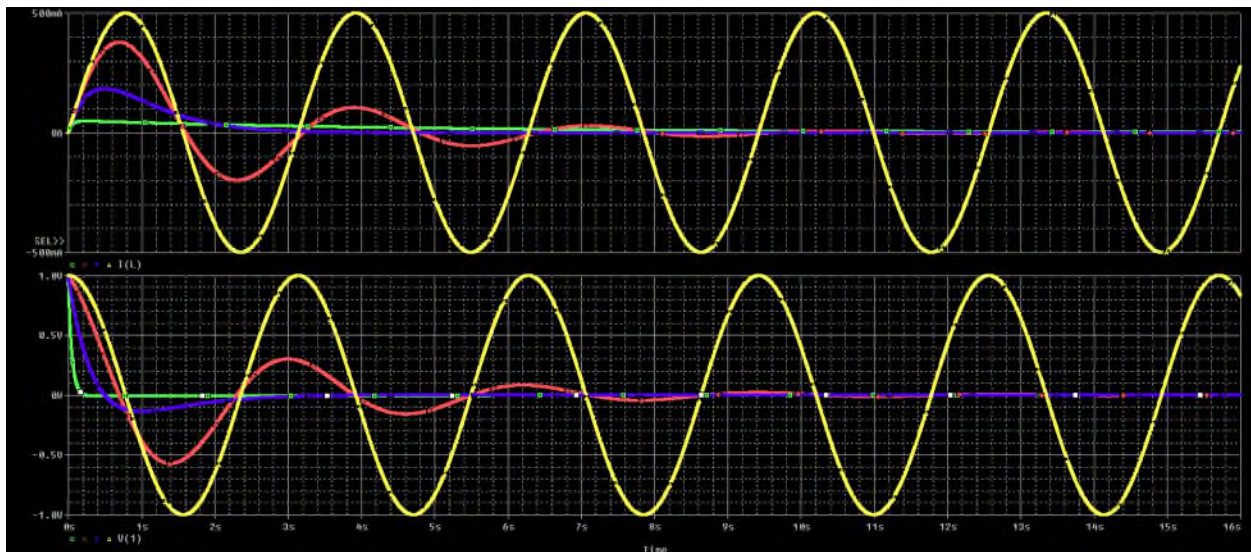


A comparison of the responses of the four parallel circuit examples (1.1, 2.1, 3.1 and 4.1) is shown below:

```

Example 1.1
C 1 0 {1/4} IC=1
G 1 0 1 0 5
L 1 0 1 IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END
Example 2.1
C 1 0 {1/4} IC=1
G 1 0 1 0 {1/5}
L 1 0 1 IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END
Example 3.1
C 1 0 {1/4} IC=1
G 1 0 1 0 1
L 1 0 1 IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END
Example 4.1
C 1 0 {1/4} IC=1
L 1 0 1 IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END

```



A comparison of the responses of the four series circuit examples (1.2, 2.2, 3.2 and 4.2) is shown below:

```

Example 1.2
L 0 1 1 IC=1
R 1 2 20
C 2 0 {1/4} IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END
Example 2.2
L 0 1 1 IC=1
R 1 2 {4/5}
C 2 0 {1/4} IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END
Example 3.2
L 0 1 1 IC=1
R 1 2 4
C 2 0 {1/4} IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END
Example 4.2
L 0 2 1 IC=1
C 2 0 {1/4} IC=0
.TRAN 1 16 0 1m UIC
.PROBE
.END

```

